EXPLICIT SOLUTION OF SECOND BVP OF THE ELASTIC MIXTURE FOR HALF-SPACE

Bitsadze L.

I. Vekua Institute of Applied Mathematics of Iv. Javakhishvili Tbilisi State University

Abstract. In this paper we consider the second BVP of elastic mixture theory for a transversally-isotropic half-space. The solution of second BVP for the transversally-isotropic half-space is given in [1]. The present paper is an attempt to extend this result to BVP of elastic mixture theory for a transversally-isotropic elastic body. Using the potential method and the theory of integral equations, the uniqueness theorem is proved for half-space and the second BVP is solved effectively (in quadratures).

Keywords and phrases: Elastic mixture, uniqueness theorem, potential method.

AMS subject classification (2000): 74E30; 74G05.

Second BVP and the uniqueness theorem for half-space

Let the plane Ox_1x_2 be the boundary of the half-space $x_3 > 0$. Let the upper halfspace will be denoted by D and the boundary of D by S. Let the axis Ox_3 be directed vertically upwards and the normal is n(0, 0, 1).

A basic equation of statics of transversally-isotropic elastic mixture theory can be written in the form [2]

$$C(\partial x)U = \begin{pmatrix} C^{(1)}(\partial x) & C^{(3)}(\partial x) \\ C^{(3)}(\partial x) & C^{(2)}(\partial x) \end{pmatrix} U = 0,$$
(1)

where

$$\begin{split} C^{(j)} &= (C^{(j)}_{pq})_{3x3}, C^{(j)}_{pq} = C^{(j)}_{qp}, j = 1, 2, 3, \\ C^{(j)}_{11}(\partial x) &= c^{(j)}_{11} \frac{\partial^2}{\partial x_1^2} + c^{(j)}_{66} \frac{\partial^2}{\partial x_2^2} + c^{(j)}_{44} \frac{\partial^2}{\partial x_3^2}, C^{(j)}_{12}(\partial x) = (c^{(j)}_{11} - c^{(j)}_{66}) \frac{\partial^2}{\partial x_1 \partial x_2}, \\ C^{(j)}_{k3}(\partial x) &= (c^{(j)}_{13} + c^{(j)}_{44}) \frac{\partial^2}{\partial x_k \partial x_3}, k = 1, 2, C^{(j)}_{22}(\partial x) = c^{(j)}_{66} \frac{\partial^2}{\partial x_1^2} + c^{(j)}_{11} \frac{\partial^2}{\partial x_2^2} \\ + c^{(j)}_{44} \frac{\partial^2}{\partial x_3^2}, C^{(j)}_{33}(\partial x) = c^{(j)}_{44} \left(\frac{\partial^2}{\partial x_1^2} + \frac{\partial^2}{\partial x_2^2} \right) + c^{(j)}_{33} \frac{\partial^2}{\partial x_3^2}, \end{split}$$

 $c_{pq}^{(k)}$ are constants characterizing the physical properties of the mixture and satisfying certain inequalities caused by the positive definiteness of potential energy. $U^T(x) = (u', u'')$ is six-dimensional displacement vector-function, $u'(x) = (u'_1, u'_2, u'_3)$ and $u''(x) = (u''_1, u''_2, u''_3)$ are partial displacement vectors. Throughout this paper the superscript "T" denotes transposition.

Definition 1. A vector-function U(x) defined in the domain D, is called regular if it has integrable continuous second derivatives in D and U(x) itself and its first derivatives are continuously extendable at every point of the boundary of D, i.e., $U(x) \in$ $C^2(D) \cap C^1(D)$ and satisfies the following conditions at infinite

$$U(x) = O(|x|^{-1}), \frac{\partial U}{\partial x_k} = O(|x|^{-2}), |x|^2 = x_1^2 + x_2^2 + x_3^2, j = 1, 2, 3; k = 1, 2, 3.$$

For the equation (1) we pose the following BVP. Find a regular function U(x), satisfying in D the equation (1), if on the boundary S the stress vector is given in the form

$$[T(\partial x, n)U]^+ = f(z), z \in S,$$
(2)

where $(.)^+$ denotes the limiting value from D and f is a given vector. $T(\partial x, n)U$ is a stress vector

$$(T(\partial x, n)U)_{k} = c_{44}^{(1)} \left(\frac{\partial u_{k}'}{\partial x_{3}} + \frac{\partial u_{3}'}{\partial x_{k}}\right) + c_{44}^{(3)} \left(\frac{\partial u_{k}''}{\partial x_{3}} + \frac{\partial u_{3}''}{\partial x_{k}}\right), k = 1, 2,$$

$$(T(\partial x, n)U)_{3} = c_{13}^{(1)} \left(\frac{\partial u_{1}'}{\partial x_{1}} + \frac{\partial u_{2}'}{\partial x_{2}}\right) + c_{13}^{(3)} \left(\frac{\partial u_{1}''}{\partial x_{1}} + \frac{\partial u_{2}''}{\partial x_{2}}\right) + c_{33}^{(1)} \frac{\partial u_{3}'}{\partial x_{3}} + c_{33}^{(3)} \frac{\partial u_{3}''}{\partial x_{3}},$$

$$(T(\partial x, n)U)_{k} = c_{44}^{(3)} \left(\frac{\partial u_{k-3}'}{\partial x_{3}} + \frac{\partial u_{3}'}{\partial x_{k-3}}\right) + c_{44}^{(2)} \left(\frac{\partial u_{k-3}''}{\partial x_{3}} + \frac{\partial u_{3}''}{\partial x_{k-3}}\right), k = 4, 5,$$

$$(T(\partial x, n)U)_{6} = c_{13}^{(3)} \left(\frac{\partial u_{1}'}{\partial x_{1}} + \frac{\partial u_{2}'}{\partial x_{2}}\right) + c_{13}^{(2)} \left(\frac{\partial u_{1}''}{\partial x_{1}} + \frac{\partial u_{3}''}{\partial x_{2}}\right) + c_{33}^{(3)} \frac{\partial u_{3}'}{\partial x_{3}} + c_{33}^{(2)} \frac{\partial u_{3}''}{\partial x_{3}}.$$

$$(3)$$

The Uniqueness Theorem. Let us prove that the second homogeneous BVP has only trivial solution. Note that, if U is the regular solution of the equation (1) and satisfies the following conditions at infinite

$$U(x) = O(|x|^{-\alpha}), P(\partial x, n)U = O(|x|^{-1-\alpha}), \alpha > 0$$

we have the following formula

$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} E(U,U) dy_1 dy_2 = -\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} U^- [TU]^- dy_1 dy_2,$$
(4)

where $T(\partial y, n)U$ is a stress vector, $E(U, U) \ge 0$. If $[TU]^- = 0$, from (4) follows U = a + [b, x], but $U(x) = O(|x|^{-\alpha})$, that a = 0, b = 0, and $U = 0, x \in D$. Therefore the homogeneous equation has only a trivial solution. Thus we shall formulate the following **Theorem.** The second BVP has at most one regular solution.

The second BVP. The solution of the second BVP will be sought in the domain D in the form

$$U(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} M(x-y)g(y)dy_1dy_2,$$
(5)

where g is an unknown real vector. M(x-y) is the following matrix

$$M(x-y) = \begin{pmatrix} \Gamma^{(1)} & \Gamma^{(2)} \\ \Gamma^{(3)} & \Gamma^{(4)} \end{pmatrix},$$
(6)

where

$$\Gamma^{(j)}(x-y) = \sum_{k=1}^{6} \|\Gamma_{pq}^{j(k)}\|_{3x3}, j = 1, 2, 3,$$

$$\Gamma^{(j)}_{pq} = \left(\delta_{pq} \frac{A_{11}^{(k)}}{r_k} + A_{12}^{(k)} \frac{\partial^2 \Phi_k}{\partial x_p \partial x_q}\right) A_k, \Gamma^{1(k)}_{3p} = A_{13}^{(k)} \frac{\partial^2 \Phi_k}{\partial x_p \partial x_3} A_k, \Gamma^{1(k)}_{p3} = A_{13}^{(k)} \frac{\partial^2 \Phi_k}{\partial x_p \partial x_3} B_k,$$

$$\Gamma^{1(k)}_{33} = \frac{A_{33}^{(k)}}{r_k} B_k, \Gamma^{2(k)}_{pq} = \left(\delta_{pq} \frac{A_{14}^{(k)}}{r_k} + A_{42}^{(k)} \frac{\partial^2 \Phi_k}{\partial x_p \partial x_q}\right) C_k, \Gamma^{2(k)}_{p3} = A_{16}^{(k)} \frac{\partial^2 \Phi_k}{\partial x_p \partial x_3} D_k,$$

$$\Gamma^{2(k)}_{3p} = A_{34}^{(k)} \frac{\partial^2 \Phi_k}{\partial x_p \partial x_3} C_k, \Gamma^{2(k)}_{33} = \frac{A_{36}^{(k)}}{r_k} D_k, \Gamma^{3(k)}_{pq} = \left(\delta_{pq} \frac{A_{14}^{(k)}}{r_k} + A_{42}^{(k)} \frac{\partial^2 \Phi_k}{\partial x_p \partial x_3}\right) A_k,$$

$$\Gamma^{3(k)}_{p3} = A_{16}^{(k)} \frac{\partial^2 \Phi_k}{\partial x_p \partial x_3} A_k, \Gamma^{3(k)}_{3p} = A_{34}^{(k)} \frac{\partial^2 \Phi_k}{\partial x_p \partial x_3} B_k, \Gamma^{3(k)}_{33} = \frac{A_{36}^{(k)}}{r_k} B_k, \Gamma^{4(k)}_{pq} = \left(\delta_{pq} \frac{A_{44}^{(k)}}{r_k} + A_{45}^{(k)} \frac{\partial^2 \Phi_k}{\partial x_p \partial x_q}\right) A_k,$$

$$+A_{45}^{(k)} \frac{\partial^2 \Phi_k}{\partial x_p \partial x_q} C_k, \Gamma^{4(k)}_{p3} = A_{46}^{(k)} \frac{\partial^2 \Phi_k}{\partial x_p \partial x_3} D_k, \Gamma^{4(k)}_{3p} = A_{46}^{(k)} \frac{\partial^2 \Phi_k}{\partial x_p \partial x_3} C_k, \Gamma^{4(k)}_{33} = \frac{A_{66}^{(k)}}{r_k} D_k,$$

$$\phi_k = x_3 \ln(x_3 + r_k) - r_k, r_k^2 = a_k [x_1 - y_1)^2 + (x_2 - y_2)^2] + x_3^2, q, p = 1, 2; \delta_{pq} = 1,$$

$$p = q, \delta_{pq} = 0, p \neq q,$$
(7)

The coefficients $A_{pq}^{(k)}$ are following

$$\begin{aligned} A_{11}^{(k)} &= \frac{(-1)^k b_0(c_{44}^{(2)} - c_{66}^{(2)} a_k)}{r_0(a_1 - a_2)}, A_{14}^{(k)} &= -\frac{(-1)^k b_0(c_{44}^{(3)} - c_{66}^{(3)} a_k)}{r_0(a_1 - a_2)}, A_{12}^{(k)} &= \frac{A_{11}^{(k)}}{a_k}, \\ A_{24}^{(k)} &= \frac{A_{14}^{(k)}}{a_k}, A_{45}^{(k)} &= \frac{A_{44}^{(k)}}{a_k}, A_{44}^{(k)} &= \frac{(-1)^k b_0(c_{44}^{(1)} - c_{66}^{(1)} a_k)}{r_0(a_1 - a_2)}, k = 1, 2, \\ A_{12}^{(k)} &= \frac{\delta_k}{a_k} [-q_3 c_{44}^{(2)} + a_k t_{12} - a_k^2 t_{11} + c_{11}^{(2)} q_4 a_k^3], A_{42}^{(k)} &= \frac{\delta_k}{a_k} [q_3 c_{43}^{(3)} + a_k t_{13} - a_k^2 t_{22} - c_{11}^{(3)} q_4 a_k^3], \\ A_{45}^{(k)} &= \frac{\delta_k}{a_k} [-q_3 c_{44}^{(1)} + a_k t_{23} - a_k^2 t_{33} + c_{11}^{(1)} q_4 a_k^3], A_{33}^{((k))} &= \delta_k [q_4 c_{33}^{(2)} - a_k t_{42} + a_k^2 t_{44} - c_{44}^{(2)} q_1 a_k^3], \\ A_{36}^{(k)} &= \delta_k [-q_4 c_{33}^{(3)} - a_k t_{62} + a_k^2 t_{66} + c_{43}^{(3)} q_1 a_k^3], A_{66}^{(k)} &= \delta_k [q_4 c_{33}^{(1)} - a_k t_{52} + a_k^2 t_{55} - c_{44}^{(1)} q_1 a_k^3], \\ A_{36}^{(k)} &= \delta_k [v_{13} - v_{11} a_k + v_{12} a_k^2], A_{16}^{(k)} &= \delta_k [w_{13} - w_{12} a_k + v_{11} a_k^2], k = 3, \dots, 6, \\ A_{34}^{(k)} &= \delta_k [v_{23} - v_{21} a_k + v_{22} a_k^2], A_{46}^{(k)} &= \delta_k [w_{34} - w_{14} a_k + w_{24} a_k^2], k = 3, \dots, 6, \\ \sum_{k=1}^6 A_{13}^{(k)} &= 0, \sum_{k=1}^6 A_{45}^{(k)} &= 0, \sum_{k=1}^6 A_{16}^{(k)} &= 0, \sum_{k=1}^6 A_{46}^{(k)} &= 0, \sum_{k=1}^6 A_{12}^{(k)} &= 0, \\ \delta_k &= d_k (a_1 - a_k) (a_2 - a_k), k = 3, \dots, 6, \end{aligned}$$

where a_k are the positive root of the characteristic equation

$$(r_0a^2 - c_0a + q_4)(b_0a^4 - b_1a^3 + b_2a^2 - b_3a + b_4) = 0,$$

$$r_0 = c_{66}^{(1)}c_{66}^{(2)} - c_{66}^{(3)2}, c_0 = c_{66}^{(1)}c_{44}^{(2)} + c_{44}^{(1)}c_{66}^{(2)} - 2c_{66}^{(3)}c_{44}^{(3)}.$$

The coefficients $d_k, b_k, v_{ij}, w_{ij}, t_{ij}$ are given in [3].

We can easily prove that every column of the matrix M(x-y) is a solution of the system (1) with respect to the point x, if $x \neq y$.

From (5) for the stress vector we obtain

$$T(\partial x, n)U(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} T(\partial x, n)M(x - y)g(y)dy_1dy_2,$$
(9)

where

$$T(\partial x, n)M(x - y) = \sum_{k=1}^{6} \begin{pmatrix} M^{(1k)} & M^{(3k)} \\ M^{(4k)} & M^{(2k)} \end{pmatrix},$$
(10)

and the elements of the matrix $M^{j(k)}(x-y), j = 1, 2, 3, 4$, can be written as

$$\begin{split} M_{pq}^{(1k)} &= \left(\delta_{pq} \gamma_{11}^{(k)} \frac{\partial}{\partial x_3} \frac{1}{r_k} + \gamma_{12}^{(k)} \frac{\partial^3 \Phi_k}{\partial x_p \partial x_q \partial x_3} \right) A_k, \\ M_{p3}^{(1k)} &= a_k \gamma_{12}^{(k)} \frac{\partial}{\partial x_p} \frac{1}{r_k} A_k, \\ M_{33}^{(1k)} &= a_k \gamma_{13}^{(k)} B_k \frac{\partial}{\partial x_3} \frac{1}{r_k}, \\ M_{33}^{(1k)} &= a_k \gamma_{12}^{(k)} \frac{\partial}{\partial x_p} \frac{1}{r_k} A_k, \\ M_{33}^{(1k)} &= a_k \gamma_{16}^{(k)} \frac{\partial}{\partial x_p} \frac{1}{r_k} A_k, \\ M_{p3}^{(3k)} &= D_k \gamma_{16}^{(k)} \frac{\partial}{\partial x_p} \frac{1}{r_k}, \\ M_{33}^{(3k)} &= D_k a_k \gamma_{16}^{(k)} B_k \frac{\partial}{\partial x_3} \frac{1}{r_k}, \\ M_{p3}^{(3k)} &= D_k a_k \gamma_{16}^{(k)} B_k \frac{\partial}{\partial x_3} \frac{1}{r_k}, \\ M_{pq}^{(4k)} &= \left(\delta_{pq} \gamma_{41}^{(k)} \frac{\partial}{\partial x_3} \frac{1}{r_k} + \gamma_{42}^{(k)} \frac{\partial}{\partial x_p} \frac{1}{r_k} A_k \right) A_k, \\ M_{p3}^{(4k)} &= B_k \gamma_{43}^{(k)} \frac{\partial}{\partial x_p} \frac{1}{r_k}, \\ M_{3p}^{(4k)} &= a_k A_k \gamma_{42}^{(k)} \frac{\partial}{\partial x_p} \frac{1}{r_k} A_{33}^{(4k)} = a_k \gamma_{43}^{(k)} B_k \frac{\partial}{\partial x_3} \frac{1}{r_k}, \\ M_{pq}^{(2k)} &= \left(\delta_{pq} \mu_{44}^{(k)} \frac{\partial}{\partial x_3} \frac{1}{r_k} + \gamma_{45}^{(k)} \frac{\partial^3 \Phi_k}{\partial x_p \partial x_q \partial x_3} \right) C_k, \\ M_{p3}^{(2k)} &= a_k \gamma_{45}^{(k)} \frac{\partial}{\partial x_p} \frac{1}{r_k} C_k, \\ M_{3p}^{(2k)} &= D_k a_k \gamma_{45}^{(k)} \frac{\partial}{\partial x_p} \frac{1}{r_k} C_k, \\ M_{3p}^{(2k)} &= A_k \gamma_{45}^{(k)} \frac{\partial}{\partial x_p} \frac{1}{r_k} C_k, \\ M_{33}^{(2k)} &= D_k a_k \gamma_{45}^{(k)} \frac{\partial}{\partial x_p} \frac{1}{r_k} C_k, \\ M_{33}^{(2k)} &= D_k a_k \gamma_{45}^{(k)} B_k \frac{\partial}{\partial x_3} \frac{1}{r_k} C_k, \\ M_{33}^{(2k)} &= D_k a_k \gamma_{45}^{(k)} B_k \frac{\partial}{\partial x_3} \frac{1}{r_k} C_k, \\ M_{32}^{(2k)} &= D_k a_k \gamma_{45}^{(k)} \frac{\partial}{\partial x_p} \frac{1}{r_k} C_k, \\ M_{33}^{(2k)} &= D_k a_k \gamma_{46}^{(k)} B_k \frac{\partial}{\partial x_3} \frac{1}{r_k} C_k, \\ M_{33}^{(2k)} &= D_k a_k \gamma_{45}^{(k)} B_k \frac{\partial}{\partial x_3} \frac{1}{r_k} C_k \\ M_{33}^{(2k)} &= D_k a_k \gamma_{45}^{(k)} \frac{\partial}{\partial x_3} \frac{1}{r_k} C_k \\ M_{33}^{(2k)} &= D_k a_k \gamma_{46}^{(k)} B_k \frac{\partial}{\partial x_3} \frac{1}{r_k} C_k \\ M_{33}^{(2k)} &= D_k a_k \gamma_{46}^{(k)} B_k \frac{\partial}{\partial x_3} \frac{1}{r_k} \\ M_{33}^{(2k)} &= D_k a_k \gamma_{46}^{(k)} B_k \frac{\partial}{\partial x_3} \frac{1}{r_k} \\ M_{33}^{(2k)} &= D_k a_k \gamma_{46}^{(k)} B_k \frac{\partial}{\partial x_3} \frac{1}{r_k} \\ M_{33}^{(2k)} &= D_k a_k \gamma_{46}^{(k)} B_k \frac{\partial}{\partial x_3} \frac{1}{r_k} \\ M_{33}^{(2k)} &= D_k a_k \gamma_{46}^{(k)} B_k \frac{\partial}{\partial x_3} \frac{1}{$$

where

$$\begin{split} \gamma_{11}^{(k)} &= c_{44}^{(1)} A_{11}^{(k)} + c_{44}^{(3)} A_{41}^{(k)}, \gamma_{12}^{(k)} = c_{44}^{(1)} (A_{12}^{(k)} + A_{13}^{(k)}) + c_{44}^{(3)} (A_{61}^{(k)} + A_{24}^{(k)}), \\ \gamma_{13}^{(k)} &= c_{44}^{(1)} (A_{13}^{(k)} + A_{33}^{(k)}) + c_{44}^{(3)} (A_{43}^{(k)} + A_{36}^{(k)}), \gamma_{14}^{(k)} = c_{44}^{(1)} A_{41}^{(k)} + c_{43}^{(3)} A_{44}^{(k)}, \\ \gamma_{16}^{(k)} &= c_{44}^{(1)} (A_{36}^{(k)} + A_{16}^{(k)}) + c_{44}^{(3)} (A_{46}^{(k)} + A_{66}^{(k)}), \gamma_{45}^{(k)} = c_{44}^{(3)} (A_{42}^{(k)} + A_{43}^{(k)}) + c_{44}^{(2)} (A_{45}^{(k)} + A_{46}^{(k)}), \\ \gamma_{43}^{(k)} &= c_{44}^{(3)} (A_{33}^{(k)} + A_{13}^{(k)}) + c_{42}^{(2)} (A_{36}^{(k)} + A_{43}^{(k)}), \gamma_{46}^{(k)} = c_{44}^{(3)} (A_{36}^{(k)} + A_{16}^{(k)}) + c_{44}^{(2)} (A_{66}^{(k)} + A_{46}^{(k)}), \\ \gamma_{24}^{(k)} &= c_{44}^{(1)} (A_{42}^{(k)} + A_{43}^{(k)}) + c_{44}^{(3)} (A_{46}^{(k)} + A_{45}^{(k)}), \gamma_{42}^{(k)} = c_{44}^{(3)} (A_{12}^{(k)} + A_{13}^{(k)}) + c_{44}^{(2)} (A_{42}^{(k)} + A_{16}^{(k)}), \\ \mu_{44}^{(k)} &= c_{44}^{(3)} A_{41}^{(k)} + c_{42}^{(2)} A_{44}^{(k)}, \gamma_{41}^{(k)} = c_{44}^{(3)} A_{11}^{(k)} + c_{41}^{(2)} A_{44}^{(k)}, k = 3, \dots, 6. \end{split}$$

For the sought for unknown constants A_k, B_k, C_k, D_k we have the following equation

$$\sum_{k=1}^{2} \frac{\gamma_{11}^{(k)}}{a_k} A_k = -1, \sum_{k=1}^{2} \frac{\gamma_{41}^{(k)}}{a_k} A_k = 0, \sum_{k=1}^{2} \frac{\gamma_{14}^{(k)}}{a_k} C_k = 0, \sum_{k=1}^{2} \frac{\mu_{44}^{(k)}}{a_k} C_k = -1,$$

$$\sum_{k=3}^{6} \gamma_{12}^{(k)} A_k = 1, \sum_{k=3}^{6} \gamma_{12}^{(k)} \sqrt{a_k} A_k = 0, \sum_{k=3}^{6} \gamma_{42}^{(k)} A_k = 0, \sum_{k=3}^{6} \gamma_{42}^{(k)} \sqrt{a_k} A_k = 0,$$

$$\sum_{k=3}^{6} \gamma_{24}^{(k)} C_k = 0, \sum_{k=3}^{6} \gamma_{24}^{(k)} \sqrt{a_k} C_k = 0, \sum_{k=3}^{6} \gamma_{45}^{(k)} C_k = 1, \sum_{k=3}^{6} \gamma_{45}^{(k)} \sqrt{a_k} C_k = 0,$$

$$\sum_{k=3}^{6} \frac{\gamma_{13}^{(k)}}{\sqrt{a_k}} B_k = 0, \sum_{k=3}^{6} \gamma_{13}^{(k)} B_k = -1, \sum_{k=3}^{6} \frac{\gamma_{43}^{(k)}}{\sqrt{a_k}} B_k = 0, \sum_{k=3}^{6} \gamma_{43}^{(k)} B_k = 0,$$

$$\sum_{k=3}^{6} \frac{\gamma_{16}^{(k)}}{\sqrt{a_k}} D_k = 0, \sum_{k=3}^{6} \gamma_{16}^{(k)} D_k = 0, \sum_{k=3}^{6} \frac{\gamma_{46}^{(k)}}{\sqrt{a_k}} D_k = 0, \sum_{k=3}^{6} \gamma_{46}^{(k)} D_k = -1.$$
(12)

By the uniqueness theorem, we conclude that the system (12) is solvable and we uniquely define A_k, B_k, C_k, D_k .

Taking into account the properties of the double layer potential and the boundary condition for determining g, from (9) we obtain the following Fredholm integral equation of second kind:

$$-g(z) + \frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} T(\partial z, n) M(x-y)g(y)dy_1dy_2 = f(z), z \in S.$$
(13)

From the last equation we have g = -f and (5) takes the form

$$U(x) = -\frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} M(x-y)f(y)dy_1dy_2, x \in D.$$
 (14)

Thus we have obtained the Poisson formula for the solution of the second BVP for the half space. Note that (14) is valid if and only if $f \in C^{0,\alpha}(S)$ and satisfies the condition $f = \frac{A}{|x|^{1+\beta}}\alpha > 0$ at infinite, where A is a constant vector and $\beta > 0$.

Acknowledgements. The designated project has been fulfilled by financial support of Georgian National Science Foundation (Grant NGNSF/ST06/3-033). Any idea in this publications is possessed by the author and may not represent the opinion of Georgian National Science Foundation itself.

REFERENCES

1. Kupradze B.D., Gegelia T.G., Basheleishvili M.O., Burchuladze T.V. Three-dimensional problems of the mathematical theory of elasticity and thermoelasticity, North-Holland Publ. Comp., Amsterdam, 1979.

2. Rushchitski Ya.Ya. Elements of Mixture Theory, (in Russian), Naukova Dumka, Kiev, 1991.

3. Bitsadze L. The basic BVPs of statics of the theory of elastic transversally isotropic mixtures for half-plane, Reports Seminar of I. Vekua Inst. of Appl. Math. **26-27**, 1-3 (2000-2001), 79-87.

Received: 12.06.2008; revised: 14.10.2008; accepted: 11.11.2008.