

THE SOLUTION OF SOME PROBLEMS OF THE THEORY OF THERMOELASTICITY WITH MICROTEmPERATURES FOR A CIRCULAR RING

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Abstract. In this work we consider the two-dimensional version of statics of the linear theory of elastic materials with inner structure whose particles, in addition to the classical displacement and temperature fields, possess microtemperatures. Some BVPs are solved for a circular ring.

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1 Basic equations. The basic system of equations of the theory of thermoelasticity with microtemperatures can be written in the form [1, 2]:

$$\begin{aligned}\mu\Delta u_\alpha + (\lambda + \mu)\partial_\alpha\theta - \beta\partial_\alpha T &= 0, \\ k_6\Delta w_\alpha + (k_4 + k_5)\partial_\alpha\Theta - k_3\partial_\alpha T - k_2w_\alpha &= 0, \\ k\Delta T + k_1\vartheta &= 0,\end{aligned}\tag{1}$$

where $\lambda, \mu, \beta, k, k_1, \dots, k_6$ are constitutive coefficients, u_α are components of the displacement vector \mathbf{u} of the point (x_1, x_2) , w_α are components of the microtemperature vector \mathbf{w} of the point (x_1, x_2) , $T(x_1, x_2)$ is the temperature measured from the constant absolute temperature T_0 ($T_0 > 0$), $\theta = \partial_1 u_1 + \partial_2 u_2$, $\Theta = \partial_1 w_1 + \partial_2 w_2$, $\Delta = \partial_{11} + \partial_{22}$ is the two dimensional Laplace operator.

On the plane Ox_1x_2 a complex variable $z = x_1 + ix_2$, where i the imaginary unit, and the following operators $\partial_z = 0.5(\partial_1 - i\partial_2)$, $\partial_{\bar{z}} = 0.5(\partial_1 + i\partial_2)$ are introduced. Then the system consisting of the equations (1) can be written in complex form as follows

$$\begin{aligned}\mu\Delta u_+ + 2(\lambda + \mu)\partial_{\bar{z}}\theta - 2\beta\partial_{\bar{z}}T &= 0, \\ k_6\Delta w_+ + 2(k_4 + k_5)\partial_{\bar{z}}\Theta - 2k_3\partial_{\bar{z}}T - k_2w_+ &= 0, \\ k\Delta T + k_1\vartheta &= 0,\end{aligned}\tag{2}$$

where $\Delta = 4\partial_z\partial_{\bar{z}}$; $u_+ := u_1 + iu_2$; $w_+ := w_1 + iw_2$.

For the positive definiteness of the corresponding quadratic form will satisfy the conditions

$$k_4 + k_5 + k_6 > 0, \quad k_2 > 0, \quad k_1k_3 - kk_2 < 0, \quad k > 0.$$

In [3] it is shown that the general solution of system (12) is represented as follows:

$$\begin{aligned}
2\mu u_+ &= \varkappa \phi(z) - z\overline{\phi'(z)} - \overline{\psi(z)} \\
&+ \frac{\mu\beta}{\lambda + 2\mu} \left\{ \frac{k_2}{2k_3} [\varphi(z) + z\overline{\varphi'(z)}] - \frac{2k_1}{kk^*} \partial_{\bar{z}} \chi_1(z, \bar{z}) \right\}, \\
w_+ &= -\overline{\varphi''(z)} + \partial_{\bar{z}} [\chi_1(z, \bar{z}) + i\chi_2(z, \bar{z})], \\
T &= \frac{k_2}{2k_3} [\varphi'(z) + \overline{\varphi'(z)}] - \frac{k_1}{2k} \chi_1(z, \bar{z}),
\end{aligned} \tag{3}$$

where $\varkappa = \frac{\lambda + 3\mu}{\lambda + \mu}$, $\varphi(z)$, $\phi(z)$ and $\psi(z)$ are the arbitrary analytic function of a complex variable z , $\chi_1(z, \bar{z})$ is a general solution of the following Helmholtz equation $\Delta \chi_1 - k^* \chi_1 = 0$, $k^* = \frac{k_2 k - k_1 k_3}{k(k_4 + k_5 + k_6)} > 0$, $\chi_2(z, \bar{z})$ is a general solution of the following Helmholtz equation $\Delta \chi_2 - \tilde{k} \chi_2 = 0$, $\tilde{k} = \frac{k_2}{k_6}$.

2 A problem for a circular ring. In this section, we solve a concrete boundary value problem for a concentric circular ring with radius R_1 and R_2 [4].

We consider the following problem

$$u_+ = \begin{cases} \sum_{n=-\infty}^{+\infty} A'_n e^{in\vartheta}, & |z| = R_1, \\ \sum_{n=-\infty}^{+\infty} A''_n e^{in\vartheta}, & |z| = R_2, \end{cases} \tag{4}$$

$$w_+ = \begin{cases} \sum_{n=-\infty}^{+\infty} B'_n e^{in\vartheta}, & |z| = R_1, \\ \sum_{n=-\infty}^{+\infty} B''_n e^{in\vartheta}, & |z| = R_2, \end{cases} \quad T = \begin{cases} \sum_{n=-\infty}^{+\infty} C'_n e^{in\vartheta}, & |z| = R_1, \\ \sum_{n=-\infty}^{+\infty} C''_n e^{in\vartheta}, & |z| = R_2. \end{cases} \tag{5}$$

The analytic functions $\phi(z)$, $\psi(z)$, $\varphi(z)$ and the metaharmonic functions $\chi_1(z, \bar{z})$, $\chi_2(z, \bar{z})$ are represented as a series

$$\phi(z) = a \ln z + \sum_{n=-\infty}^{+\infty} a_n z^n, \quad \psi(z) = b + \sum_{n=-\infty}^{+\infty} b_n z^n, \quad \varphi'(z) = c \ln z + \sum_{n=-\infty}^{+\infty} c_n z^n, \tag{6}$$

$$\chi_\gamma(z, \bar{z}) = \sum_{n=-\infty}^{+\infty} \left(\alpha_{\gamma n} I_n(\sqrt{k^*} r) + \beta_{\gamma n} K_n(\sqrt{k^*} r) \right) e^{in\vartheta}, \tag{7}$$

where $I_n(\cdot)$ and $K_n(\cdot)$ are modified Bessel functions of order n .

In the boundary conditions (5) we substitute the corresponding expressions (6), (7) and compare the coefficients at identical degrees. We obtain the following system of

equations

$$\begin{aligned} \frac{k_2}{k_3} c \ln R_\gamma + \frac{k_2}{2k_3} c_0 - \frac{k_1}{2k} \left[\alpha_{10} I_0(\sqrt{k^*} R_\gamma) + \beta_{10} K_0(\sqrt{k^*} R_\gamma) \right] &= C_{\gamma 0}, \\ -\frac{2c}{R_\gamma} + \sqrt{k^*} \left[\alpha_{10} I_1(\sqrt{k^*} R_\gamma) - \beta_{10} K_1(\sqrt{k^*} R_\gamma) \right] &= B_{\gamma 1} + \bar{B}_{\gamma 1}, \\ \sqrt{\tilde{k}} \left[\alpha_{20} I_1(\sqrt{\tilde{k}} R_\gamma) - \beta_{20} K_1(\sqrt{\tilde{k}} R_\gamma) \right] &= -i(B_{\gamma 1} - \bar{B}_{\gamma 1}), \end{aligned} \quad (8)$$

$$\begin{aligned} \frac{k_2}{k_3} \left[R_\gamma^n c_n + R_\gamma^{-n} \bar{c}_{-n} \right] - \frac{k_1}{2k} \left[\alpha_{1n} I_n(\sqrt{k^*} R_\gamma) + \beta_{1n} K_n(\sqrt{k^*} R_\gamma) \right] &= C_{\gamma n}, \\ -n \left[R_\gamma^{n-1} c_n - R_\gamma^{-n-1} \bar{c}_{-n} \right] & \\ + \frac{\sqrt{k^*}}{2} \left[(I_{n-1}(\sqrt{k^*} R_\gamma) + I_{n+1}(\sqrt{k^*} R_\gamma)) \alpha_{1n} - (K_{n-1}(\sqrt{k^*} R_\gamma) + K_{n+1}(\sqrt{k^*} R_\gamma)) \beta_{1n} \right] & \\ -i \frac{\sqrt{\tilde{k}}}{2} \left[(I_{n-1}(\sqrt{\tilde{k}} R_\gamma) - I_{n+1}(\sqrt{\tilde{k}} R_\gamma)) \alpha_{2n} - (K_{n-1}(\sqrt{\tilde{k}} R_\gamma) - K_{n+1}(\sqrt{\tilde{k}} R_\gamma)) \beta_{2n} \right] & \\ = B_{\gamma n+1} + \bar{B}_{\gamma 1-n}, & \\ -n \left[R_\gamma^{n-1} c_n + R_\gamma^{-n-1} \bar{c}_{-n} \right] & \\ + \frac{\sqrt{k^*}}{2} \left[(I_{n-1}(\sqrt{k^*} R_\gamma) - I_{n+1}(\sqrt{k^*} R_\gamma)) \alpha_{1n} - (K_{n-1}(\sqrt{k^*} R_\gamma) - K_{n+1}(\sqrt{k^*} R_\gamma)) \beta_{1n} \right] & \\ -i \frac{\sqrt{\tilde{k}}}{2} \left[(I_{n-1}(\sqrt{\tilde{k}} R_\gamma) + I_{n+1}(\sqrt{\tilde{k}} R_\gamma)) \alpha_{2n} - (K_{n-1}(\sqrt{\tilde{k}} R_\gamma) + K_{n+1}(\sqrt{\tilde{k}} R_\gamma)) \beta_{2n} \right] & \\ = -B_{\gamma n+1} + \bar{B}_{\gamma 1-n}, \quad \gamma = 1, 2. & \end{aligned} \quad (9)$$

It is possible to show that the determinant of the systems (8) and (9) are other than zero. Equations (8), (9) uniquely determine all coefficients of decomposition of functions $\varphi(z)$, $\chi_1(z, \bar{z})$ and $\chi_2(z, \bar{z})$. Absolute and uniform convergence of the series obtained in the ring (including the contours) will be provided if the functions set on the boundary have sufficient smoothness.

In the boundary conditions (5) we substitute the corresponding expressions (6), (7):

$$\begin{aligned} (\varkappa a + \bar{b}) i \vartheta + (\varkappa a - \bar{b}) \ln R_\gamma - \bar{a} e^{2i\vartheta} + \sum_{-\infty}^{+\infty} \varkappa a_n R_\gamma^n e^{in\vartheta} & \\ - \sum_{-\infty}^{+\infty} (n \bar{a}_n R_\gamma^n e^{-i(n-2)\vartheta} + \bar{b}_n R_\gamma^n e^{-in\vartheta}) &= E i \vartheta + \sum_{-\infty}^{+\infty} D_\gamma n e^{in\vartheta}, \quad \gamma = 1, 2, \end{aligned} \quad (10)$$

where

$$\begin{aligned} E &= -\frac{\mu\beta}{\lambda + 2\mu} \frac{k_2}{2k_3} c_{-1}, \\ D_{\gamma 0} &= -\frac{\mu\beta}{\lambda + 2\mu} \left\{ \frac{k_2}{2k_3} [\ln R_\gamma + 1] c_{-1} - \frac{k_1}{k\sqrt{k^*}} [I_0(\sqrt{k^*} R_\gamma) \alpha_{1-1} - K_0(\sqrt{k^*} R_\gamma) \beta_{1-1}] \right\}, \\ D_{\gamma 1} &= -\frac{\mu\beta}{\lambda + 2\mu} \frac{k_2}{2k_3} [c R_\gamma (2 \ln R_\gamma - 1) + 2c_0 R_\gamma] \\ &+ \frac{\mu\beta}{\lambda + 2\mu} \frac{k_1}{k\sqrt{k^*}} \left[I_1(\sqrt{k^*} R_\gamma) \alpha_{10} - K_1(\sqrt{k^*} R_\gamma) \beta_{10} \right], \end{aligned}$$

$$D_{\gamma n} = -\frac{\mu\beta}{\lambda + 2\mu} \frac{k_2}{2k_3} \left[\frac{c_{n-1}}{n} R_{\gamma}^n + \bar{c}_{1-n} R_{\gamma}^{-n} \right] \\ + \frac{\mu\beta}{\lambda + 2\mu} \frac{k_1}{k\sqrt{k^*}} \left[I_n(\sqrt{k^*} R_{\gamma}) \alpha_{1n-1} - K_n(\sqrt{k^*} R_{\gamma}) \beta_{1n-1} \right].$$

The requirement for uniqueness of displacements have the following form

$$\varkappa a + \bar{b} = E. \quad (11)$$

Comparison of terms involving $e^{in\vartheta}$ gives

$$\begin{aligned} \varkappa \ln R_{\gamma} a - 2R_{\gamma}^2 \bar{a}_2 + \varkappa a_0 - \bar{b}_0 &= D_{\gamma 0}, \\ -\bar{a} + \varkappa R_{\gamma}^2 a_2 - R_{\gamma}^{-2} \bar{b}_{-2} &= D_{\gamma 2}, \\ \varkappa R_{\gamma}^n a_n + (n-2) R_{\gamma}^{-n+2} \bar{a}_{-n+2} - R_{\gamma}^{-n} \bar{b}_{-n} &= D_{\gamma n}, \quad \gamma = 1, 2. \end{aligned} \quad (12)$$

Equations (11), (12) uniquely determine all coefficients of decomposition of functions $\phi(z)$ and $\psi(z)$.

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