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## THE PROBLEM OF EXISTENCE THE NEUTRAL SURFACE FOR THE ELASTIC SHELL $^{\star}$

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**Abstract**. I. Vekua obtained the conditions for the existence of the neutral surface of a shell, when the neutral surface is the middle surface. In this paper the neutral surface is considered as any equidistant surfaces of the shell.

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The stress-strain relations are given in the form [1, 2]

$$\sigma_{j}^{i} = \lambda \theta g_{j}^{i} + 2\mu e_{j}^{i}, \quad (i, j = 1, 2, 3), \tag{1}$$

where  $\sigma_j^i$  and  $e_j^i$  are the mixed components, respectively of stress and strain tensors,  $\theta$  is the cubical dilatation which will be written as

$$\theta = \theta' + e_3^3, \quad \theta' = e_\alpha^\alpha, \quad (\alpha = 1, 2). \tag{2}$$

when j = 3 from (1) we have

$$\sigma_3^{\alpha} = 2\mu e_3^{\alpha}, \ \ \sigma_3^3 = \lambda\theta + 2\mu e_3^3 = \lambda\theta' + (\lambda + 2\mu)e_3^3.$$
 (3)

From (3)

$$e_3^{\alpha} = \frac{1}{2\mu}\sigma_3^{\alpha}, \quad e_3^3 = -\frac{\lambda}{\lambda+2\mu}\theta' + \frac{1}{\lambda+2\mu}\sigma_3^3. \tag{4}$$

By inserting (4) into (2) we obtain

$$\theta = \frac{\lambda'}{\lambda}\theta' + \frac{1}{\lambda + 2\mu}\sigma_3^3, \quad \lambda' = \frac{2\lambda\mu}{\lambda + 2\mu}.$$
(5)

Substituting expression (5) into (1) we get

$$\sigma_j^i = T_j^i + Q_j^i = \left(\lambda'\theta' + \frac{\lambda}{\lambda + 2\mu}\sigma_3^3\right)g_j^i + 2\mu e_j^i,$$

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where

$$T^{\alpha}_{\beta} = \lambda' \theta g^{\alpha}_{\beta} + 2\mu e^{\alpha}_{\beta}, \ Q^{\alpha}_{\beta} = \sigma' \sigma^3_3 g^{\alpha}_{\beta}, \ T^i_3 = 0, \ Q^i_3 = \sigma^i_3, \ \left(\sigma' = \frac{\lambda}{\lambda + 2\mu}\right).$$
(6)

The vector  $\mathbf{T}^{\alpha}$  satisfies the condition  $\mathbf{nT}^{\alpha} = 0$  and is therefore called the tangential stress field and the vector  $\mathbf{Q}^{i}$  will be called the transverse field.

The vectorial equation of equilibrium

$$\frac{1}{\sqrt{g}}\partial_i(\sqrt{g}\boldsymbol{\sigma}^i) + \boldsymbol{\Phi} = 0, \quad (\sqrt{g} = \sqrt{a}\vartheta, \quad \vartheta = 1 - 2Hx_3 + Kx_3^2)$$
(7)

may be written as

$$\frac{1}{\sqrt{g}} [\partial_{\alpha} (\sqrt{g} \boldsymbol{T}^{\alpha}) + \partial_{i} (\sqrt{g} \boldsymbol{Q}^{i})] + \boldsymbol{\Phi} = 0.$$
(8)

Let the surface  $\hat{S}$ :  $x^3 = \text{const}$  be the neutral surface of a non-shallow shell. Then  $T^{\alpha} = 0$ , i.e.  $T^{\alpha\beta} = 0$  (on  $\hat{S}$ ), and equation (8)

$$\frac{1}{\sqrt{a}}\partial_{\alpha}(\sqrt{a}\vartheta\boldsymbol{Q}^{\alpha})+\partial_{3}(\vartheta\boldsymbol{\sigma}^{3})+\vartheta\boldsymbol{\Phi}=0,$$

or

$$\left[\nabla_{\alpha}(\vartheta \boldsymbol{Q}^{\alpha}) + \partial_{3}(\vartheta \boldsymbol{\sigma}^{3}) + \vartheta \boldsymbol{\Phi}\right]_{x^{3} = \text{const}} = 0, \quad (-h \le x^{3} = x_{3} \le h)$$
(9)

where 2h is the thickness of the shell and

$$\boldsymbol{Q}^{\alpha} = \sigma' \sigma_3^3 \boldsymbol{R}^{\alpha} + \sigma_3^{\alpha} \boldsymbol{n}, \quad \nabla_{\alpha}(\cdot) = \frac{1}{\sqrt{a}} \partial_{\alpha}(\sqrt{a} \cdot). \tag{10}$$

Denote the stress forces acting on the face surfaces  $S^+$  and  $S^-$  by  $\stackrel{(+)}{P}$  and  $\stackrel{(-)}{P}$ . We have

$$\overset{(+)}{P} = -(\boldsymbol{\sigma}^3)_{x^3=h}, \quad \overset{(-)}{P} = (\boldsymbol{\sigma}^3)_{x^3=-h}.$$
(11)

If we approximately represent  $\sigma^3$  by the formula

$$\boldsymbol{\sigma}^{3}(x^{1}, x^{2}, x^{3}) \cong \overset{0}{\boldsymbol{\sigma}}(x^{1}, x^{2}) + x^{3} \overset{1}{\boldsymbol{\sigma}}(x^{1}, x^{2}).$$
(12)

from (11) we get

$$\boldsymbol{\sigma}^{3}(x^{1}, x^{2}, x^{3}) \cong -\frac{1}{2} \Big[ \stackrel{(+)}{\boldsymbol{P}} - \stackrel{(-)}{\boldsymbol{P}} + \frac{x^{3}}{h} \Big( \stackrel{(+)}{\boldsymbol{P}} + \stackrel{(-)}{\boldsymbol{P}} \Big) \Big] \\ = -\frac{1}{2} \Big[ \frac{h + x_{3}}{h} \Big( \stackrel{(+)}{\boldsymbol{P}} - \stackrel{(-)}{\boldsymbol{P}} \Big) + \frac{2x^{3}}{h} \stackrel{(-)}{\boldsymbol{P}} \Big],$$
(13)

or

$$\boldsymbol{\sigma}^{3}(x^{1}, x^{2}, x^{3}) = -\frac{1}{2} \Big[ \frac{h + x_{3}}{h} \Big( P^{\alpha} \mathbf{r}_{\alpha} + P^{3} \mathbf{n} \Big) + \frac{2x^{3}}{h} \frac{(-)}{P} \Big], \tag{14}$$

$$P^{\alpha} = P^{\alpha} - P^{\alpha}, \quad P = P^{3} - P^{3},$$

Then to define the vector field  $\stackrel{(+)}{P}$  we have the equation

$$\left\{\nabla_{\alpha}(\sigma'A^{\alpha}_{\beta}P\boldsymbol{r}^{\beta} + AP^{\alpha}\boldsymbol{n}) + B(P\boldsymbol{n} + P^{\alpha}\boldsymbol{r}_{\alpha}) + \tilde{\boldsymbol{\Phi}}\right\}_{x^{3}=c} = 0$$
(15)

where

$$A^{\alpha}_{\beta} = \frac{h+c}{h} [a^{\alpha}_{\beta} + c(b^{\alpha}_{\beta} - 2Ha^{\alpha}_{\beta})], \quad A = \frac{h+c}{h} \vartheta(c),$$
  

$$B = \frac{1}{h} [1 - 2Hh + 2(Kh - 2H)c + 3Kc^{2}],$$
  

$$\tilde{\Phi} = -2\vartheta(c)\Phi(c) + \nabla_{\alpha} \left\{ \sigma' \frac{2c}{h} \left[ a^{\alpha}_{\beta} + c(b^{\alpha}_{\beta} - 2Ha^{\alpha}_{\beta}) \right] \overset{(-)}{P^{3}} \boldsymbol{r}^{\beta} + \frac{2c}{h} \vartheta(c) \overset{(-)}{P^{\alpha}} \boldsymbol{n} \right\}$$
  

$$+ \frac{2}{h} \left[ \vartheta(c) + 2(Kc - H) \right] \overset{(-)}{\boldsymbol{P}}.$$
(16)

Since

$$abla_{lpha}oldsymbol{r}^{eta}=b_{lpha}^{eta}oldsymbol{n}, \ \ ext{and} \ \ 
abla_{lpha}oldsymbol{n}=-b_{lphaeta}oldsymbol{r}^{eta}$$

from (15) we have

$$\sigma' \nabla_{\alpha} (A^{\alpha}_{\beta} P) + (Ba_{\alpha\beta} - Ab_{\alpha\beta}) P^{\alpha} + \tilde{\Phi}_{\beta} = 0, \quad (\tilde{\Phi}_{\beta} = \tilde{\Phi} \boldsymbol{r}_{\beta}), \tag{17}$$

$$\nabla_{\alpha}(AP^{\alpha}) + (\sigma' A^{\alpha}_{\beta} b^{\beta}_{\alpha} + B)P + \tilde{\Phi}_{3} = 0, \quad (\tilde{\Phi}_{3} = \tilde{\Phi}\boldsymbol{n}).$$
(18)

From the system of equations (17) we have

$$P^{\alpha} = \overset{(+)}{P^{\alpha}} - \overset{(-)}{P^{\alpha}} = -\hat{d}^{\alpha\beta} \left[ \nabla_{\alpha} (A^{\alpha}_{\beta} P) + \tilde{\Phi}_{\beta} \right], \qquad (19)$$

where

$$\hat{d}^{\alpha\beta} = \frac{1}{\Delta} \left[ (B - 2AH) a^{\alpha\beta} + Ab^{\alpha\beta} \right], \quad \hat{F}_{\beta} = - \left[ \tilde{\Phi}_{\beta} + \nabla_{\alpha} (A^{\alpha}_{\beta} P) \right], \quad (20)$$
$$\Delta = B^2 - 2ABH + A^2 K.$$

Inserting expressions (19) into (18) we obtain the equation

$$\sigma' \nabla_{\alpha} \left[ A \hat{d}^{\alpha\beta} \nabla_{\gamma} (A^{\alpha}_{\beta} P) \right] - (B + \sigma' A^{\alpha}_{\beta} b^{\beta}_{\alpha}) P + \Phi = 0.$$
<sup>(21)</sup>

It is easily seen that equation (21) is of the elliptic type.

Thus, if the surface  $x^3 = c$  is neutral, the stress  $\stackrel{(+)}{P}$  and  $\stackrel{(+)}{P}$ , applied to the face surfaces, must satisfy the vector equation (17) and (18). This means that the stresses  $\stackrel{(+)}{P}$  and  $\stackrel{(-)}{P}$ cannot be prescribed arbitrarily both at the same time. However there are problems when this does not occur. For example, in aircraft or submarine apparatus the force  $\stackrel{(-)}{P}$  acting on the inner surface  $S^-$  may be assumed to be prescribed, but the  $\stackrel{(T)}{P}$  acting on the external face surface  $S^+$  is not, in general, assigned beforehand. The same situation occurs on dams. One face surface of the dam is free from stresses and the other is under the hydrodynamic load, a variable which is generally difficult to define exactly at any moment in time.

## $\mathbf{R} \to \mathbf{F} \to \mathbf{R} \to \mathbf{N} \to \mathbf{C} \to \mathbf{S}$

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