

A DIFFERENCE SCHEME REPRESENTATION FOR A NONLINEAR
KIRCHHOFF EQUATION

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Abstract. The initial boundary value problem for the dynamic string equation $w_{tt} - \left(\lambda + \frac{2}{\pi} \int_0^L w_x^2(x, t) dx\right) w_{xx}(x, t) = 0$ is considered. To solve it, the difference scheme is written, which is represented in the form convenient for both solution and investigation since the eigenfunctions of a difference operator are used as a basis.

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1. Statement of the problem. Let us consider the nonlinear equation

$$w_{tt}(x, t) - \left(\lambda + \frac{2}{\pi} \int_0^\pi w_x^2(x, t) dx\right) w_{xx}(x, t) = 0, \quad (1)$$

$$0 < x < \pi, \quad 0 < t \leq T,$$

with the initial boundary conditions

$$w(x, 0) = w^0(x), \quad w_t(x, 0) = w^1(x), \quad (2)$$

$$w(0, t) = w(\pi, t) = 0, \quad (3)$$

$$0 \leq x \leq \pi, \quad 0 \leq t \leq T.$$

Equation (1) describing the string vibration was obtained by Kirchhoff in 1876 [1]. A great number of works is dedicated to the investigation of this equation and its generalizations (see e.g. [2] and the bibliography therein).

2. Difference scheme. On the intervals $[0, \pi]$ and $[0, T]$ of the change of arguments x and t we introduce the nets $\omega_h = \{x_i = ih, i = 0, 1, \dots, N\}$ and $\omega_\tau = \{t_m = m\tau, m = 0, 1, \dots, M\}$, and to the rectangle $[0, \pi] \times [0, T]$ we put into correspondence the net

$$\Omega_{h\tau} = \omega_h \times \omega_\tau = \{(x_i, t_m), i = 0, 1, \dots, N, m = 0, 1, \dots, M\},$$

where h and τ are the steps for which we have $h = \frac{\pi}{N}$, $\tau = \frac{T}{M}$. The value of some function defined on the net $\Omega_{h\tau}$ at the node (x_i, t_m) is denoted by w_i^m . Let us approximate equation (1) and conditions (2), (3) by means of a difference scheme which using the standard notation [3] is written in the form

$$w_{tt,i}^m - \frac{1}{2} \left(\lambda + \frac{1}{\pi} \sum_{p=1,-1}^N h \sum_{j=1}^N (w_{\bar{x},j}^{m+p})^2 \right) \sum_{r=1,-1} w_{\bar{x},i}^{m+r} = 0, \quad (4)$$

$$i = 1, 2, \dots, N-1, \quad m = 1, 2, \dots, M-1,$$

$$w_i^0 = \omega_i^0, \quad w_i^1 = \omega_i^1, \quad i = 1, 2, \dots, N-1, \quad (5)$$

$$w_0^m = w_N^m = 0, \quad m = 0, 1, \dots, M, \quad (6)$$

where $\omega_i^0 = w_i^0$, $\omega_i^1 = w^0(x_i) + \tau w^1(x_i) + \frac{\tau^2}{2} \left(\lambda + \frac{2}{\pi} \int_0^\pi (w^{0r}(x))^2 dx \right) w^{0r}(x_i)$.

Under the error of the difference scheme (4)-(6) we understand the net function

$$\begin{aligned} \Delta w_i^m &= w(x_i, t_m) - w_i^m. \\ i &= 0, 1, \dots, N, \quad m = 0, 1, \dots, M. \end{aligned} \quad (7)$$

A system of equations for the error has the form

$$\begin{aligned} \Delta w_{tt,i}^m &- \frac{1}{2} \left(\lambda + \frac{1}{\pi} \sum_{p=1,-1} h \sum_{j=1}^N (w_{\bar{x},j}^{m+p})^2 \right) \sum_{r=1,-1} \Delta w_{\bar{x},i}^{m+r} \\ &- \frac{1}{2\pi} \left[\left(\sum_{p=1,-1} h \sum_{j=1}^N (w_{\bar{x}}(x_j, t_{m+p}) + w_{\bar{x},j}^{m+p}) \right) \Delta w_{\bar{x},j}^{m+p} \right] \sum_{r=1,-1} w_{\bar{x},i}^{m+r} \\ &= \psi_i^{m+1,m-1}, \quad i = 1, 2, \dots, N-1, \quad m = 1, 2, \dots, M-1, \end{aligned} \quad (8)$$

where

$$\begin{aligned} \psi_i^{m+1,m-1} &= w_{tt}(x_i, t_m) - \frac{1}{2} \left(\lambda + \frac{1}{\pi} \sum_{p=1,-1} h \sum_{j=1}^N (w_{\bar{x}}(x_j, t_{m+p}))^2 \right) \\ &\quad \times \sum_{r=1,-1} w_{\bar{x}}(x_i, t_{m+r}). \end{aligned}$$

From (5),(6) and (2), (3) we obtain

$$\Delta w_i^0 = 0, \quad \Delta w_i^1 = \frac{\tau^3}{3!} w_{ttt}(x_i, \theta_i), \quad 0 \leq \theta_i \leq \tau, \quad i = 1, 2, \dots, N-1, \quad (9)$$

$$\Delta w_0^m = \Delta w_N^m = 0, \quad m = 0, 1, \dots, M. \quad (10)$$

3. Replacement of the basis. Using the values of the solution of the difference scheme (4)-(6) at the internal nodes of the set ω_h on the m -th layer, i.e. for $t = t_m$, we form the vector \mathbf{w}^m . Thus $\mathbf{w}^m = (w_i^m)_{i=1}^{N-1}$. Let us write the vector \mathbf{w}^m in terms of the basis $\{\mathbf{e}^i\}_{i=1}^{N-1}$, where the basis vector \mathbf{e}^i is the ort $\mathbf{e}^i = (\delta_{ij})_{j=1}^{N-1}$, δ_{ij} is the Kronecker symbol. We have

$$\mathbf{w}^m = \sum_{i=1}^{N-1} w_i^m \mathbf{e}^i. \quad (11)$$

Let us replace the basis. On the net ω_h , we consider the following problem of eigenvalues: find a net function μ_j , $j = 0, 1, \dots, N$, and a constant λ such that

$$\mu_{\bar{x},j} + \lambda \mu_j = 0, \quad j = 1, 2, \dots, N-1, \quad (12)$$

$$\mu_0 = \mu_N = 0. \quad (13)$$

As it is known, at any rate for sufficiently large N there exist $N-1$ linearly independent solutions of this problem, the i -th solution, $i = 1, 2, \dots, N-1$, has the form [3] $\mu_j^i =$

$\sqrt{\frac{2h}{\pi}} \sin ijh$, $\lambda_i = \frac{4}{h^2} \sin^2 \frac{ih}{2}$, $j = 1, 2, \dots, N-1$, $\mu_0^i = \mu_N^i = 0$, also, $A = (A_{ij})_{i,j=1}^{N-1}$, where $A_{ij} = \mu_j^i$ is an orthonormalized matrix.

Let us consider the basis $\{\boldsymbol{\mu}^i\}_{i=1}^{N-1}$, where the basis vector $\boldsymbol{\mu}^i = (\mu_j^i)_{j=1}^{N-1}$. We write the vector \boldsymbol{w}^m in the new basis

$$\boldsymbol{w}^m = \sum_{i=1}^{N-1} v_i^m \boldsymbol{\mu}^i, \quad m = 0, 1, \dots, M. \quad (14)$$

Let us complement the set of coefficients of expansion (14) with values $v_0^m = v_N^m = 0$ and rewrite the difference scheme (4)-(6) using v_i^m . Taking (11)-(14) into account, we obtain

$$v_{tt,i}^m + \frac{1}{2} \lambda_i \left(\lambda + \frac{1}{\pi} \sum_{p=1,-1} h \sum_{j=1}^{N-1} \lambda_j (v_j^{m+p})^2 \right) \sum_{r=1,-1} v_i^{m+r} = 0, \quad (15)$$

$$i = 1, 2, \dots, N-1, \quad m = 1, 2, \dots, M-1,$$

$$v_i^0 = \nu_i^0, \quad v_i^1 = \nu_i^1, \quad i = 1, 2, \dots, N-1, \quad (16)$$

$$v_0^m = v_N^m = 0, \quad m = 0, 1, \dots, M. \quad (17)$$

Here $\nu_1^p, \nu_2^p, \dots, \nu_{N-1}^p$ is the solution of the system of linear algebraic equations $\sum_{j=1}^{N-1} A_{ij} \nu_j = \omega_i^p$, $i = 1, 2, \dots, N-1$, $p = 0, 1$.

Now, let rewrite the system for error (8)-(10) in the new basis. Using the values Δw_i^m defined by (7) we construct the vector $\Delta \boldsymbol{w}^m = (\Delta w_i^m)_{i=1}^{N-1}$ and write it in the form

$$\Delta \boldsymbol{w}^m = \sum_{i=1}^{N-1} \Delta v_i^m \boldsymbol{\mu}^i, \quad m = 0, 1, \dots, M. \quad (18)$$

Let us complement the set of coefficients of expansion (18) with values $\Delta v_0^m = \Delta v_N^m = 0$. We introduce into consideration the vector $\boldsymbol{w}(x, t_m) = (w(x_i, t_m))_{i=1}^{N-1}$ and the values $v(x_i, t_m)$ from the expansion

$$\boldsymbol{w}(x, t_m) = \sum_{i=1}^{N-1} v(x_i, t_m) \boldsymbol{\mu}^i. \quad (19)$$

Taking into account (12)-(14) and (18), (19), by (8)-(10) we obtain

$$\begin{aligned} & \Delta v_{tt,i}^m + \frac{1}{2} \lambda_i \left(\lambda + \frac{1}{\pi} \sum_{p=1,-1} h \sum_{j=1}^{N-1} \lambda_j (v_j^{m+p})^2 \right) \sum_{r=1,-1} \Delta v_i^{m+r} \\ & + \frac{1}{2\pi} \lambda_i \left(\sum_{p=1,-1} h \sum_{j=1}^{N-1} \lambda_j (v(x_j, t_{m+p}) + v_j^{m+p}) \Delta v_j^{m+p} \right) \sum_{r=1,-1} v_i^{m+r} \\ & = \varphi_i^{m+1, m-1}, \quad i = 1, 2, \dots, N-1, \quad m = 1, 2, \dots, M-1, \end{aligned} \quad (20)$$

$$\Delta v_i^0 = 0, \quad \Delta v_i^1 = \Delta \omega_i, \quad i = 1, 2, \dots, N-1, \quad (21)$$

$$\Delta v_0^m = \Delta v_N^m = 0, \quad m = 0, 1, \dots, M. \quad (22)$$

Here $\varphi_1^{m+1,m-1}, \varphi_2^{m+1,m-1}, \dots, \varphi_{N-1}^{m+1,m-1}$ and $\Delta\omega_1, \Delta\omega_2, \dots, \Delta\omega_{N-1}$ are respectively the solutions of the systems of algebraic equations $\sum_{j=1}^{N-1} A_{ij}\varphi_j = \psi_i^{m+1,m-1}, i = 1, 2, \dots, N-1$, and $\sum_{j=1}^{N-1} A_{ij}\Delta\omega_j = \Delta w_i^1, i = 1, 2, \dots, N-1$.

Comparing the initial and the transformed systems, we come to a conclusion that for solution of the difference scheme, system (15)–(17) is more convenient than system (4)–(6), whereas in investigating the difference method convergence, it is less difficult to obtain a priori estimates from system (20)–(22) as compared with the case of using system (8)–(10). To conclude, it should be noted that by solving the system layer-by-layer and finding $v_1^m, v_2^m, \dots, v_{N-1}^m$ from (15)–(17), we obtain, by virtue of (14), $w_1^m, w_2^m, \dots, w_{N-1}^m$.

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R E F E R E N C E S

1. Kirchhoff G. Vorlesungen über Mathematische Physik. I. Mechanik. (German) *Teubner, Leipzig*, 1876.
2. Peradze J. An approximate algorithm for a Kirchhoff wave equation. *SIAM J. Numer. Anal.* **47**, 3 (2009), 2243–2268.
3. Samarskii A.A. Theory of Difference Schemes. (Russian) *Second edition. Nauka, Moscow*, 1983.

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