Reports of Enlarged Session of the Seminar of I. Vekua Institute of Applied Mathematics Volume 25, 2011

ON ONE NONLINEAR DIFFUSION SYSTEM

Gagoshidze M., Jangveladze T.

Abstract. One-dimensional nonlinear diffusion system of Maxwell's equations by taking into account Joule's rule and thermal conductivity is considered. Finite difference schemes and splitting-up models are constructed. Graphs of respective numerical experiments are given.

Keywords and phrases: One-dimensional nonlinear Maxwell's system, Joule's rule, thermal conductivity, finite difference schemes, splitting-up models.

AMS subject classification: 35Q60, 35Q61, 35K55, 65M06, 83C50.

System of Maxwell's equations describing process of penetration of magnetic field into a substance by taking into account Joule's rule and thermal conductivity, has the following form [1]:

$$\frac{\partial \theta}{\partial t} = \nu_m (rot H)^2 + div(\kappa \operatorname{grad} \theta), \quad \frac{\partial H}{\partial t} = -rot \ (\nu_m rot H), \tag{1}$$

where $H = (H_1, H_2, H_3)$ is a vector of the magnetic field, θ is temperature, ν_m and κ are characteristics coefficients of the substance. As a rule these coefficients are functions of argument θ .

Let us consider following initial-boundary value problem for one-dimensional analog of the system (1) with special power-type nonlinearity:

$$\frac{\partial U}{\partial t} = \frac{\partial}{\partial x} \left(V^{\alpha} \frac{\partial U}{\partial x} \right), \quad \frac{\partial V}{\partial t} = V^{\alpha} \left(\frac{\partial U}{\partial x} \right)^{2} + \frac{\partial^{2} V}{\partial x^{2}}, \qquad (x,t) \in \Omega \times (0,T), \\
U(x,t) = \frac{\partial V(x,t)}{\partial x} = 0, \qquad (x,t) \in \partial\Omega \times (0,T), \\
U(x,0) = U_{0}(x), \quad V(x,0) = V_{0}(x) \ge Const > 0, \qquad x \in \overline{\Omega},$$
(2)

where $-1/2 \leq \alpha \leq 1/2$; U_0 , V_0 are known functions defined on $\overline{\Omega} = [0, 1]$, T is the fixed positive constant.

After introducing new unknown function $V^{1/2} = W$ and notation $2\alpha = \gamma$ the problem (2) can be reduced to the following equivalent form:

$$\frac{\partial U}{\partial t} = \frac{\partial}{\partial x} \left(W^{\gamma} \frac{\partial U}{\partial x} \right), \qquad (x,t) \in \Omega \times (0,T),
\frac{\partial W}{\partial t} = \frac{1}{2} W^{\gamma-1} \left(\frac{\partial U}{\partial x} \right)^{2} + \frac{\partial^{2} W}{\partial x^{2}} + \frac{1}{W} \left(\frac{\partial W}{\partial x} \right)^{2}, \qquad (x,t) \in \Omega \times (0,T),
U(x,t) = \frac{\partial W(x,t)}{\partial x} = 0, \qquad (x,t) \in \partial\Omega \times (0,T),
U(x,0) = U_{0}(x), \quad W(x,0) = W_{0}(x) = V_{0}^{1/2}(x), \qquad x \in \overline{\Omega},$$
(3)

where $-1 \leq \gamma \leq 1$.

Let us note that same reduction for the problem (2) is also used in [2] as well as for investigation of problem analogical to (2) with first kind of boundary conditions for function V in the work [3]. The questions of existence, uniqueness, regularity of the solutions and numerical resolution of the initial-boundary value problems to the (1), (2) and (3) type models and related parabolic systems are discussed in many works (see, for example, [4]-[20] and references therein).

Note that, system (1) without thermal conductivity can be reduced to integrodifferential form [21]. Many responses were followed after publication of the paper [21] (see, for example, [22]-[28]). The questions of existence, uniqueness, asymptotic behavior of the solutions and numerical resolution of some kind of initial-boundary value problems for this type integro-differential models are studied in these works.

It is important to construct and study discrete analogs for the model (1). It is also necessary to investigate this question for one-dimensional analog (2).

Complex nonlinearity dictates also to split the investigated model along the physical process and to investigate basic model by them (see, for example, [2], [3]).

Different type of splitting-up schemes are constructed and investigated for many models of mathematical physics (see, for example, [29]-[31] and references therein).

Investigation of splitting-up schemes along the physical processes for one-dimensional analog of system (1) is the natural beginning of studding this issue. In this direction investigations was made in the works [2], [3] and in a number of other works as well. In the paper [3] semi-discrete additive models for the initial-boundary value problem for system (2) with Dirichlet boundary conditions on temperature is considered. The same question for the second kind of boundary conditions on the function V is considered in [2]. We should note that in [2] the above mentioned splitting-up scheme is constructed for the multidimensional system (1) as well.

The aim of this note is to construct finite-difference schemes for the problem (2). The splitting-up schemes for the problem (3) are considered as well. Let us note that finite difference schemes for (2) type problems are constructed in many works (see, for example, [3], [8], [9], [16], [18] and references therein).

On $[0,1] \times [0,T]$ let us introduce a net with mesh points denoted by $(x_i, t_j) = (ih, j\tau)$, where i = 0, 1, ..., M; j = 0, 1, ..., N with h = 1/M, $\tau = T/N$. The initial line is denoted by j = 0. The discrete approximation at (x_i, t_j) is designed by u_i^j, v_i^j and the exact solution to the problem (2) by U_i^j, V_i^j .

Let us consider problem (2) with nonhomogeneous right hand sides in both equations of the system and construct the following difference scheme:

$$\frac{u_i^{j+1} - u_i^j}{\tau} = \frac{1}{h^2} \left\{ a \left(v_i^{j+1} \right) u_{i-1}^{j+1} - \left(a \left(v_i^{j+1} \right) + a \left(v_{i+1}^{j+1} \right) \right) u_i^{j+1} + a \left(v_{i+1}^{j+1} \right) u_{i+1}^{j+1} \right\} + f_i^{j+1},$$

$$(4)$$

$$\frac{v_i^{j+1} - v_i^j}{\tau} = \sigma a \left(v_i^{j+1} \right) \left(\frac{u_i^{j+1} - u_{i-1}^{j+1}}{h^2} \right)^2 + (1 - \sigma) a \left(v_i^j \right) \left(\frac{u_i^j - u_{i-1}^j}{h} \right)^2 + \frac{v_{i+1}^{j+1} - 2v_i^{j+1} + v_{i-1}^{j+1}}{h^2} + g_i^{j+1},$$
(5)

$$i = 1, 2...M - 1; j = 0, 1..., N - 1,$$

$$u_0^j = u_M^j = v_{x,0}^j = v_{\bar{x},M}^j = 0, \quad j = 0, 1, ..., N,$$
 (6)

$$u_j^0 = U_{0,i} \quad v_j^0 = V_{0,i}, \quad i = 0, 1..., M.$$
 (7)

Here $a(v) = v^{\alpha}$, $\alpha \in R$; f, g are given functions and $0 \leq \sigma \leq 1$ is also a given constant.

The difference scheme (4)-(7) is the first order, i.e. its rate is $O(\tau + h)$.

Theorem 1. If $-1/2 \leq \alpha \leq 1/2$, $0 < \sigma \leq 1$ and problem (2) has a sufficiently smooth solution, then the difference scheme (4)-(7) is stable, uniquely solvable and its solution converges to the solution of the problem (2) as $\tau \to 0, h \to 0$. The following estimate is also true

$$\left\| U^{i} - u^{j} \right\|_{h} + \left\| V^{j} + v^{j} \right\|_{h} = O(\tau + h).$$

Here $\|\cdot\|_h$ is a discrete analog of the norm of the space $L_2(0,1)$.

Note that, in the case $\sigma = 0$, for solving the finite difference scheme (4)-(7) at first we solve system (5) by known tridiagonal matrix algorithm and after we solve system (4) by same algorithm, using in both cases suitable boundary and initial conditions from (6), (7). In the case $\sigma \neq 0$, i.e. for solving fully implicit scheme (4)-(7) we must include a method of solving system of nonlinear algebraic equations. So, it is necessary to use Newton iterative process. It is easy to notice that for the solution of the implicit scheme (4)-(7) the numerical algorithms [3] based on a modified Newton method without main changes can be used.

Now, let us consider the splitting-up model for the problem (3). Using the notations:

$$y_t = \frac{y^{j+1} - y^j}{\tau}, \quad \eta_1 + \eta_2 = 1, \quad \eta_1 > 0, \quad \eta_2 > 0,$$

$$y = \eta_1 y_1 + \eta_2 y_2, \quad y_{1t} = \frac{y_1^{j+1} - y^j}{\tau}, \quad y_{2t} = \frac{y_2^{j+1} - y^j}{\tau}$$

let us correspond to the initial-boundary value problem (3) the following additive averaged semi-discrete scheme [2], [3]:

$$u_{1t} = \frac{d}{dx} \left(w_1^{\gamma} \frac{du_1}{dx} \right), \quad \eta_1 w_{1t} = \frac{1}{2} w_1^{\gamma - 1} \left(\frac{du_1}{dx} \right)^2, \quad u_{2t} = \frac{d}{dx} \left(w_2^{\gamma} \frac{du_2}{dx} \right), \\ \eta_2 w_{2t} = \frac{d^2 w_2}{dx^2} + \frac{1}{w_2} \left(\frac{dw_2}{dx} \right)^2, \quad u_1^0 = u_2^0 = U_0, \quad \frac{dw_1^0}{dx} = \frac{dw_2^0}{dx} = W_0.$$
(8)

The following statement takes place [2].

Theorem 2. If problem (3) has a sufficiently smooth solution, then the solution of scheme (8) converges to the solution of problem (3) as $\tau \to 0$ and the following estimate is true $\|U(t_j) - u^j\| + \|W(t_j) + w^j\| = O(\tau^{1/2}).$

Here $\|\cdot\|$ is an usual norm of the space $L_2(0,1)$.

As we already mentioned the additive models analogical to (8) for the multidimensional system (1) are also constructed and investigated (see, for example, [2]).

At the end we will note that numerical test experiments carried out on the basis on difference schemes (4)-(7) are founded on the above described algorithms (see Fig. 1).

Let us also note that fully discrete analog based on the semi-discrete splitting model (8) is constructed and used. Various numerical experiments are done using these algorithms. These experiments fully agree with theoretical results (see Fig. 2).

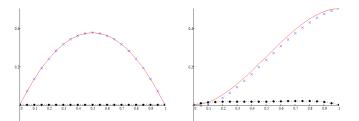


Fig. 1. Exact (solid line) and numerical (marked with \times) solutions and differences between exact and numerical solutions (marked with \bullet).

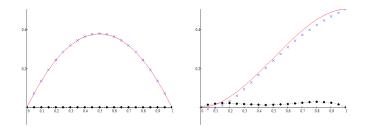


Fig. 2. Exact (solid line) and numerical (marked with \times) solutions and differences between exact and numerical solutions (marked with \bullet).

REFERENCES

1. Landau L., Lifschitz E. Electrodynamics of Continuous Media. (Russian) Moscow, 1958.

2. Jangveladze T. Additive models for one nonlinear diffusion system. *Rep. Enlarged Sess. Semin. I. Vekua Appl. Math.*, **24** (2010), 24-30.

3. Abuladze I.O., Gordeziani D.G., Dzhangveladze T.A., Korshia T.K. On the numerical modeling of a nonlinear problem of the diffusion of a magnetic field with regard to heat conductivity. (Russian) *Proc. I. Vekua Inst. Appl. Math.*, **18** (1986), 48-67.

4. Friedman A. Partial Differential Equations of Parabolic Type. Prentice-Hall, 1964.

5. Ladyzhenskaya O.A., Solonnikov V.A. Ural'ceva N.N. Linear and Quasilinear Equations of Parabolic Type. (Russian) *Moscow*, 1968.

6. Lions J.-L., Quelques Mthodes de Resolution des Problems aux Limites Non Linaires. Dunod, Gauthier-Villars, (French) Paris, 1969.

7. Dafermos C.M., Hsiao L. Adiabatic shearing of incompressible fluids with temperature dependent viscosity. *Quart. Appl. Math.*, **41**, 1 (1983), 45-58.

8. Dzhangveladze T. The difference scheme for one system of nonlinear partial differential equations. (Russian) *Rep. Enlarged Sess. Semin. I.Vekua Appl. Math.*, **2**, 3 (1986), 40-43.

9. Dzhangveladze T.A. On the convergence of the difference scheme for one nonlinear system of partial differential equations. (Russian) Bull. Acad. Sci. Georgian SSR, **126**, 2 (1987), 257-260.

10. Dzhangveladze T.A. The boundary value problem for one system of nonlinear partial differential equations. (Russian) Abstracts of All-Union School. Functional Meth. Appl. Math. Math. Phys., Tashkent, I (1988), 22-23.

11. Dzhangveladze T.A. A system of nonlinear partial differential equation. (Russian) Rep. Enlarged Sess. Semin. I. Vekua Appl. Math., 4, 1 (1989), 38-41.

12. Cimatti G. Existence of weak solutions for the nonstationary problem of the Joule heating of a conductor. Ann. Mat. Pura Appl., 162, 4 (1992), 33-42.

13. Yin,H.-M. Global solutions of Maxwell's equations in an electromagnetic field with a temperaturedependent electrical conductivity. *European J. Appl. Math.*, **5**, 1 (1994), 57-64.

14. Elliott C.M., Larsson S. A finite element model for the time-dependent Joule heating problem. *Math. Comp.*, **64** (1995), 1433-1453.

15. Bien M. Existence of global weak solutions for a class of quasilinear equations describing Joule's heating. *Math. Meth. Appl. Sci.*, **23** (1998), 1275-1291.

16. Kiguradze Z. The difference scheme for one system of nonlinear partial differential equations. *Rep. Enlarged Sess. Semin. I. Vekua Appl. Math.*, **14**, 3 (1999), 67-70.

17. Yin H.-M. On a nonlinear Maxwell's system in quasi-stationary electromagnetic fields. *Math. Models Methods Appl. Sci.*, **14**, 10 (2004), 1521-1539.

18. Sun D., Manoranjan V.S., Yin H.-M. Numerical solutions for a coupled parabolic equations arising induction heating processes. *Discrete Contin. Dyn. Syst., Supplement*, (2007), 956-964.

19. Ding S.J., Guo B.L., Lin J.Y., Zeng M. Global existence of weak solutions for Landau– Lifshitz-Maxwell equations. *Discrete Contin. Dyn. Syst.*, **17**, 4 (2007), 867-890.

20. Baxevanis Th., Katsaounis Th., Tzavaras A.E. Adaptive finite element computations of shear band formation. *Math. Mod. Meth. Appl. Sci.*, **20**, 3 (2010), 423-448.

21. Gordeziani D.G., Dzhangveladze T.A., Korshia T.K. Existence and uniqueness of the solution of a class of nonlinear parabolic problems. *Differ. Uravn.*, **19**, 7 (1983), 1197-1207 (Russian). English translation: *Differ. Equ.*, **19**, 7 (1984), 887-895.

22. Dzhangveladze T.A. First boundary-value problem for a nonlinear equation of parabolic type. (Russian) *Dokl. Akad. Nauk SSSR*, **269**, 4 (1983), 839-842. English translation: *Soviet Phys. Dokl.*, **28**, 4 (1983), 323-324.

23. Laptev G.I. Mathematical singularities of a problem on the penetration of a magnetic field into a substance. (Russian) Zh. Vychisl. Mat. Mat. Fiz., 28 (1988), 1332-1345. English translation: U.S.S.R. Comput. Math. Math. Phys., 28 (1990), 35-45.

24. Lin Y., Yin H.-M. Nonlinear parabolic equations with nonlinear functionals. J. Math. Anal. Appl., 168, 1 (1992), 28-41.

25. Jangveladze T. Convergence of a difference scheme for a nonlinear integro-differential equation. *Proc. I. Vekua Inst. Appl. Math.*, **48** (1998), 38-43.

26. Jangveladze T., Kiguradze Z., Neta B. Finite difference approximation of a nonlinear integrodifferential system. *Appl. Math.Comput.*, **215**, 2 (2009), 615-628.

27. Kiguradze Z. On asymptotic behavior and numerical resolution of one nonlinear Maxwell's model. *Proceedings of the 15th WSEAS Int. Conf. Applied Math. (MATH '10)*, (2010), 55-60.

28. Jangveladze T. Investigation and numerical solution of system of nonlinear integro-differential equations associated with the penetration of a magnetic field in a substance. *Proceedings of the 15th WSEAS Int. Conf. Applied Math.(MATH '10)*, (2010), 79-84.

29. Janenko N.N. The Method of Fractional Steps for Multi-dimensional Problems of Mathematical Physics. (Russian) *Moscow*, 1967.

30. Marchuk G.I. The Splitting-up Methods. (Russian) Moscow, 1988.

31. Samarskii A,A., Vabishchevich P.N. Additive Schemes for Mathematical Physics Problems. (Russian) *Moscow*, 1999.

Received 24.04.2011; accepted 10.09.2011.

Authors' addresses:

T. Jangveladze

I. Vekua Institute of Applied MathematicsIv. Javakhishvili Tbilisi State University2, University St., Tbilisi 0186

Caucasus University Kostava Av. 77, 0175 Georgia E-mail: tjangv@yahoo.com

M. Gagoshidze Iv. Javakhishvili Tbilisi State University 2, University St., Tbilisi 0186 Georgia E-mail: MishaGagoshidze@gmail.com