

ON A DIFFERENCE SCHEME FOR SOLUTION OF SYSTEM OF EQUATIONS
OF DYNAMICAL MULTIDIMENSIONAL PROBLEMS OF THE ELASTICITY
THEORY AND SHELL THEORY

Komurjishvili O., Khomeriki N.

Iv. Javakhishvili Tbilisi State University
N. Muskhelishvili Institute of Computational Mathematics

At first we consider the problem for the equations of elasticity theory. Let $\overline{Q}_T = \overline{G} \times [0, T]$, $Q_T = G \times (0, T)$, where $\overline{G} = \{0 \leq x_j \leq \ell_\alpha, \alpha = 1, 2, \dots, p\}$, be the p -dimensional parallelepiped.

We consider the problem of finding of solution for the system

$$\frac{\partial^2 \mathbf{u}}{\partial t^2} = L\mathbf{u} + \mathbf{f}(\mathbf{x}, \mathbf{t}), \quad (\mathbf{x}, \mathbf{t}) \in \mathbf{Q}_T, \quad (1)$$

which satisfied the following complementary conditions

$$\begin{aligned} \mathbf{u} &= \mathbf{0}, \text{ if } \mathbf{x} \in \Gamma = \partial\mathbf{G}, \mathbf{t} \in [0, \mathbf{T}], \\ \mathbf{u}(\mathbf{x}, 0) &= \mathbf{u}_0(\mathbf{x}), \quad \frac{\partial \mathbf{u}(\mathbf{x}, 0)}{\partial \mathbf{t}} = \overline{\mathbf{u}}_0(\mathbf{x}) \text{ if } \mathbf{x} \in \overline{\mathbf{G}}, \end{aligned} \quad (2)$$

where

$$L\mathbf{u} = \mu \Delta \mathbf{u} + (\lambda + \mu) \text{graddiv} \mathbf{u}, \quad \Delta \mathbf{u} = \sum_{\alpha=1}^p \frac{\partial^2 \mathbf{u}}{\partial \mathbf{x}_\alpha^2},$$

$\lambda = \text{const} > 0$, $\mu = \text{const} > 0$ are the Lamé coefficients, $\mathbf{u} = (\mathbf{u}_1, \mathbf{u}_2, \dots, \mathbf{u}_n)$ is n -dimensional vector-function. The condition of positive definiteness is fulfilled for $c_1 = \mu$, $c_2 = \lambda + 2\mu$ (see [1]).

Let $\overline{\omega}_h = \{x_i = (i_1 h_1, \dots, i_p h_p)\}$ be a net in \overline{G} , $0 \leq i_j \leq N_\alpha$, $h_\alpha = \ell_\alpha / N_\alpha$, $\alpha = 1, 2, \dots, p$ and $\overline{\omega}_\tau = \{t_j = j\tau, j = 0, 1, \dots\}$ be a net on the interval $0 \leq t \leq T$.

The set of the points (t_j, x_i) , $x_i \in \overline{\omega}_h(\omega_h)$ and $t_j \in \omega_\tau$ is denoted by $\overline{G}_{h\tau}(G_{h\tau})$.

For net functions and their difference derivatives which are defined on $\overline{G}_{h\tau}$, we take the following notations: $y^{j+1}(x) = y(t_{j+1}, x_i) = y((j+1)\tau, x_1, \dots, x_p)$, $\hat{y} = y(t_{j+2}, x_i)$, $\check{y} = y(t_j, x_i)$, $y_{\bar{t}} = (y - \check{y})/\tau$, $y_t = (\hat{y} - y)/\tau$, $y^{(-1)} = y(t_{j+1}, x_1, \dots, x_i - h_i, \dots, x_p)$, $y^{(+1)} = y(t_{j+1}, x_1, \dots, x_i + h_i, \dots, x_p)$, $y_{\bar{x}_i} = (y - y^{(-1)})/h_i$, $y_{x_i} = (y^{(+1)} - y)/h_i$, $y_{x_i}^\circ = (y^{(+1)} - y^{(-1)})/2h_i = \frac{1}{2}(y_{x_i} + y_{\bar{x}_i})$.

Let us introduce a space H of net functions defined on ω_h and vanishing on Γ_h with the inner product $(\mathbf{y}, \mathbf{v}) = \sum_{\mathbf{i}=1}^{\mathbf{n}} (\mathbf{y}_i, \mathbf{v}_i)$,

$$(y_i, v_i) = \sum_{x \in \omega_h} y_i(x) v_i(x) h_1 \dots h_p = \sum_{i_1=1}^{N_1-1} \sum_{i_2=1}^{N_2-1} \dots \sum_{i_p=1}^{N_p-1} y_i(i_1 h_1, \dots, i_p h_p)$$

$$\cdot v_i(i_1 h_1, \dots, i_p h_p) h_1 \dots h_p$$

and the norm $\|\mathbf{y}\| = \sqrt{(\mathbf{y}, \mathbf{y})}$. Besides we introduce the notations $A = A_1 + \dots + A_p$, $A_i = -\Delta_{ii}$, $\Delta_{ii}\mathbf{y} = \mathbf{y}_{\mathbf{x}_i \bar{\mathbf{x}}_i}$.

The operator A is selfadjoint and positive in H . The norm in the energetic space H_A has the form

$$\|\mathbf{y}\|_A^2 = \sum_{i_1=1}^{N_1} \sum_{i_2=1}^{N_2-1} \dots \sum_{i_p=1}^{N_p-1} \mathbf{y}_{\bar{\mathbf{x}}_1}^2 \mathbf{h}_1 \dots \mathbf{h}_p + \dots + \sum_{i_1=1}^{N_1-1} \sum_{i_2=1}^{N_2-1} \dots \sum_{i_p=1}^{N_p} \mathbf{y}_{\bar{\mathbf{x}}_p}^2 \mathbf{h}_1 \dots \mathbf{h}_p$$

or

$$\|\mathbf{y}\|_A^2 = \sum_{i=1}^p \|\mathbf{y}_{\bar{\mathbf{x}}_i}\|_i^2.$$

By analogy with [2], we put to the problem (1), (2) in the contrespondence the difference problem

$$(E + \sigma\tau^2 R_j^i) \mathbf{y}_{\mathbf{j}\mathbf{t}\bar{\mathbf{t}}} = \mathbf{L}_{\mathbf{j}\mathbf{h}}^i \mathbf{y} + \mathbf{f}_{\mathbf{j}}, \quad \mathbf{j} = \mathbf{1}, \mathbf{2}, \dots, \mathbf{n}, \quad (3)$$

$$\mathbf{y} = \mathbf{0} \quad \text{if } \mathbf{x} \in \Gamma_{\mathbf{h}}, \quad \mathbf{t} \in \bar{\omega}_{\tau}, \quad \mathbf{y}(\mathbf{x}, \mathbf{0}) = \mathbf{u}_0(\mathbf{x}), \quad \mathbf{y}_{\tau}(\mathbf{x}, \mathbf{0}) = \tilde{\mathbf{u}}_0(\mathbf{x}) \quad \text{if } \mathbf{x} \in \bar{\omega}_{\mathbf{h}}, \quad (4)$$

where

$$L_{\mathbf{j}\mathbf{h}}^i \mathbf{y} = \sum_{\mathbf{j}=1}^p \mathbf{a}_{\mathbf{j}\mathbf{j}}^i \mathbf{y}_{\mathbf{i}\mathbf{x}_i \bar{\mathbf{x}}_j} + 2\mathbf{a}_{\alpha\beta} \sum_{\substack{\alpha, \beta=1 \\ \alpha \neq \beta}}^p \Delta_{\alpha\beta} \mathbf{y},$$

$$\Delta_{\alpha\beta} \mathbf{y} = \frac{1}{4}(\mathbf{y}_{\mathbf{x}_{\alpha}} + \mathbf{y}_{\bar{\mathbf{x}}_{\alpha}})_{\mathbf{x}_{\beta}} + \frac{1}{4}(\mathbf{y}_{\mathbf{x}_{\alpha}} + \mathbf{y}_{\bar{\mathbf{x}}_{\alpha}})_{\bar{\mathbf{x}}_{\beta}}, \quad \alpha \neq \beta,$$

$$R_j^i = \frac{1}{2} a_{jj}^i A_j, \quad a_{jj}^i = \lambda + 2\mu, \quad a_{jj}^i = \mu, \quad i \neq j, \quad a_{\alpha\beta} = \frac{1}{2}(\lambda + \mu), \quad i = 1, 2, \dots, p.$$

For example we write out the absolute approximating difference scheme for $p = 2$

$$(E + \sigma\tau^2 R_j^i) \mathbf{y}_{\mathbf{j}\mathbf{t}\bar{\mathbf{t}}} = \mathbf{L}_{\mathbf{j}\mathbf{h}}^i \mathbf{y} + \mathbf{f}_{\mathbf{j}}, \quad \mathbf{i}, \mathbf{j} = \mathbf{1}, \mathbf{2},$$

where

$$L_{1h}^1 \mathbf{y} = \mathbf{a}_{11}^1 \mathbf{y}_{1\mathbf{x}\bar{\mathbf{x}}_1} + 2\mathbf{a}_{12} \Delta_{12} \mathbf{y}_2 + \mathbf{a}_{22}^1 \mathbf{y}_{1\mathbf{x}_2 \bar{\mathbf{x}}_2}, \quad L_{2h}^2 \mathbf{y} = \mathbf{a}_{22}^2 \mathbf{y}_{2\mathbf{x}\bar{\mathbf{x}}_2} + 2\mathbf{a}_{12} \Delta_{12} \mathbf{y}_1 + \mathbf{a}_{11}^2 \mathbf{y}_{2\mathbf{x}_1 \bar{\mathbf{x}}_1},$$

$$\mathbf{a}_{11}^1 = \mathbf{a}_{22}^2 = \lambda + 2\mu, \quad a_{11}^1 = a_{22}^2 = \mu, \quad a_{12} = \frac{1}{2}(\lambda + \mu), \quad \Delta_{12} \mathbf{y} = \frac{1}{4}(\mathbf{y}_{\mathbf{x}_1} + \mathbf{y}_{\bar{\mathbf{x}}_1})_{\mathbf{x}_2} + \frac{1}{4}(\mathbf{y}_{\mathbf{x}_1} + \mathbf{y}_{\bar{\mathbf{x}}_1})_{\bar{\mathbf{x}}_2},$$

$$R_j^1 = \frac{1}{2} a_{jj}^1 A_j, \quad R_j^2 = \frac{1}{2} a_{jj}^2 A_j, \quad j = 1, 2.$$

For the difference schemes of a variant of shell theory (see [3]) we obtain a system of approximating equations of order N with $3N + 3$ unknown functions. If we temporarily don't take into consideration the terms M_j^k (see [3], ch. I, §10) which contain the unknown functions and their first order derivatives, then the resulting system of equations is subdivided into two groups: the first group contains the system of equations of plane problem of elasticity theory and the second — the Poisson equation of two space variables. Therefore, for them we write out the difference schemes analogously.

Further, the derivatives of first order are approximated by the central differences

$$y_{x_i}^o.$$

The investigation of the received finite difference schemes is based on the general regularization principle (see [1], ch. VI, §3). The stability sufficient condition of the scheme (3) is given in the form

$$D > \frac{1}{4}\tau^2 L_{jh}^i \quad \text{or} \quad D > \frac{1+\varepsilon}{4}\tau^2 L_{jh}^i, \quad (5)$$

where $D = E + \sigma\tau^2 R_j^i$ and L_{jh}^i are selfadjoint positive operators. Before to receiving of the stability sufficient conditions of the scheme (3), we write out some inequalities. Taking into account the conditions of ellipticity of operator L , it is possible to show that

$$c_1 \|\mathbf{y}\|_{\mathbf{A}}^2 \leq (\mathbf{L}_{jh}^i \mathbf{y}, \mathbf{y}) \leq c_2 \|\mathbf{y}\|_{\mathbf{A}}^2.$$

Farther $\delta_1 E < A < \delta_2 E$, where E is identity operator

$$\delta_1 = \sum_{i=1}^p \frac{4}{h_i^2} \sin^2 \frac{\pi h_i}{2\ell_i}, \quad \delta_2 = \sum_{i=1}^p \frac{4}{h_i^2} \cos^2 \frac{\pi h_i}{2\ell_i}.$$

Hence we obtain $c_1 \delta_1 E \leq L_{jh}^i \leq c_2 \delta_2 E$. Then, it is possible make to use the inequalities

$$D > \frac{1}{4}\tau^2 c_2 A \quad \text{or} \quad D \geq \frac{1}{4}\tau^2 c_2 \delta_2 E,$$

instead (5). From this we obtain the following sufficient condition of stability:

$$\frac{\tau}{h} \leq \frac{1}{\sqrt{c_2(p-2\sigma)}}, \quad 0 < \sigma \leq \frac{1}{2},$$

$$\frac{\tau}{h} < \frac{1}{\sqrt{c_2(p-1)}}, \quad \sigma > \frac{1}{2},$$

$$(h = h_1 = \dots = h_p, \quad \ell_i = 1, \quad i = 1, 2, \dots, p).$$

Remark. In the interval $\frac{1}{\sqrt{c_2 p}} < \frac{\tau}{h} \leq \frac{1}{\sqrt{c_2(p-2\sigma)}}$, $\left(0 < \sigma \leq \frac{1}{2}\right)$, where the stability conditions of explicit scheme are not satisfied in the family of scheme sets with the approximation $O(\tau^2 + |h|^2)$, the represented types of schemes have preferences compared with other economical schemes.

R E F E R E N C E S

1. Samarski A.A. Theory of difference schemes, M.: Nauka, 1977 (in Russian).
2. Komurjishvili O.P. Three-layer difference schemes for multidimensional hyperbolic equation, Seminar of I.Vekua institute of Applied Mathematics, Reports, vol. 18, No. 1.
3. Vekua I.N. The theory of sloping shells with variable thickness. Tbilisi, Mecniereba, 1965 (in Russian).

Received 7.06.2006; revised 28.02.2007; accepted 25.12.2007.