

I. VEKUA APPROXIMATION $N = 1$ FOR THE NON-SHALLOW
CYLINDRICAL SHELLS

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We shall consider the non-shallow cylindrical shells for I.N. Vekua approximation $N = 1$. The displacement vector $\mathbf{U}(x^1, x^2, x^3)$ are expressed by the following formula [1]

$$\mathbf{U}(x^1, x^2, x^3) = \mathbf{u}(x^1, x^2) + \frac{x^3}{h} \mathbf{v}(x^1, x^2).$$

Here $\mathbf{u}(x^1, x^2)$ and $\mathbf{v}(x^1, x^2)$ are the vector fields on the middle surface $x^3 = 0$, $2h$ is the thickness of the shell, x^3 is a thickness coordinate ($-h \leq x^3 \leq h$), x^1 and x^2 are isometric coordinates on the cylindrical surface.

Let us construct the solutions of the form [2], [3]:

$$\begin{aligned} u_3 &= \sum_{k=0}^{\infty} u_3^k \varepsilon^k, & u_\alpha &= \sum_{k=0}^{\infty} u_\alpha^k \varepsilon^k \quad (\alpha = 1, 2), \\ v_3 &= \sum_{k=0}^{\infty} v_3^k \varepsilon^k, & v_\alpha &= \sum_{k=0}^{\infty} v_\alpha^k \varepsilon^k \quad (\alpha = 1, 2), \end{aligned}$$

where u_1, u_2 and v_1, v_2 are the tangent components of the vectors \mathbf{u} and \mathbf{v} respectively, u_3 and v_3 denote the normal components of these vectors, $\varepsilon = \frac{h}{R_0}$ is a small parameter, R_0 -the radius of the middle surface of the cylinder.

The system of equilibrium equations of the two-dimensional non-shallow cylindrical shells may be written in the following form [4], [5]:

$$\begin{aligned} \mu \Delta u_1^k &+ (\lambda + \mu) \partial_1 \theta^k + \lambda \partial_1 v_3^k = X_1^k, \\ \mu \Delta u_2^k &+ (\lambda + \mu) \partial_2 \theta^k + \lambda \partial_2 v_3^k = X_2^k, \\ \mu \Delta v_3^k &- 3 \left[\lambda \theta^k + (\lambda + 2\mu) v_3^k \right] = X_3^k, \end{aligned} \tag{1}$$

$$\begin{aligned} \mu \Delta v_1^k &+ (\lambda + \mu) \partial_1 \Theta^k - 3\mu (\partial_1 u_3^k + v_1^k) = X_4^k, \\ \mu \Delta v_2^k &+ (\lambda + \mu) \partial_2 \Theta^k - 3\mu (\partial_2 u_3^k + v_2^k) = X_5^k, \\ \mu \Delta u_3^k &+ \mu \Theta^k = X_6^k, \\ &(k = 0, 1, 2, \dots), \end{aligned} \tag{2}$$

where

$$\begin{aligned}
 X_1^k &= X_1 - \mu \left(v_1^{k-1} + \frac{\partial_{22} v_1^{k-1}}{3} \right) + \sum_{s=0}^{\lfloor \frac{k-1}{2} \rfloor} \left(\frac{\lambda + 2\mu}{2s+3} \partial_{11} v_1^{k-2s-1} - \frac{\lambda + 3\mu}{2s+1} \partial_1 u_3^{k-2s-1} \right) - \\
 &\quad - \mu \sum_{s=1}^{\lfloor \frac{k-1}{2} \rfloor} \frac{v_1^{k-2s-1}}{2s+1} - \sum_{s=1}^{\lfloor \frac{k}{2} \rfloor} \left(\frac{\lambda + 2\mu}{2s+1} \partial_{11} u_1^{k-2s} - \frac{\lambda + 3\mu}{2s+1} \partial_1 v_3^{k-2s} - \frac{\mu}{2s-1} u_1^{k-2s} \right), \\
 X_2^k &= X_2 - \lambda \partial_2 u_3^{k-1} - \frac{\lambda + 2\mu}{3} \partial_{22} v_2^{k-1} - \mu \left(\sum_{s=1}^{\lfloor \frac{k}{2} \rfloor} \frac{\partial_{11} u_2^{k-2s}}{2s+1} - \sum_{s=0}^{\lfloor \frac{k-1}{2} \rfloor} \frac{\partial_{11} v_2^{k-2s-1}}{2s+3} \right), \\
 X_3^k &= X_3 + 3\lambda u_3^{k-1} + (2\lambda - \mu) \partial_2 v_2^{k-1} - \mu \partial_{22} u_3^{k-1} + 3 \sum_{s=0}^{\lfloor \frac{k-1}{2} \rfloor} \left(\frac{\mu}{2s+3} \partial_{11} u_3^{k-2s-1} + \right. \\
 &\quad \left. + \frac{\lambda + 3\mu}{2s+3} \partial_1 v_1^{k-2s-1} \right) - 3(\lambda + 2\mu) \sum_{s=1}^{\lfloor \frac{k-1}{2} \rfloor} \frac{u_3^{k-2s-1}}{2s+1} - \\
 &\quad - 3 \sum_{s=1}^{\lfloor \frac{k}{2} \rfloor} \left(\frac{\mu}{2s+1} \partial_{11} v_3^{k-2s} + \frac{\lambda + 3\mu}{2s+1} \partial_1 u_1^{k-2s} - \frac{\lambda + 2\mu}{2s-1} v_3^{k-2s} \right), \\
 X_4^k &= X_4 - \mu \partial_{22} u_1^{k-1} + 3 \sum_{s=0}^{\lfloor \frac{k-1}{2} \rfloor} \left(\frac{\lambda + 2\mu}{2s+3} \partial_{11} u_1^{k-2s-1} - \frac{\lambda + 3\mu}{2s+3} \partial_1 v_3^{k-2s-1} + \right. \\
 &\quad \left. + \frac{\mu}{2s+1} u_1^{k-2s-1} \right) - 3 \sum_{s=1}^{\lfloor \frac{k}{2} \rfloor} \left(\frac{\lambda + 2\mu}{2s+3} \partial_{11} v_1^{k-2s} - \frac{\lambda + 3\mu}{2s+1} \partial_1 u_3^{k-2s} - \frac{\mu}{2s+1} v_1^{k-2s} \right), \\
 X_5^k &= X_5 - (\lambda + 2\mu) \partial_{22} u_2^{k-1} - (2\lambda - \mu) \partial_2 v_3^{k-1} - \\
 &\quad - 3\mu \left(\sum_{s=1}^{\lfloor \frac{k}{2} \rfloor} \frac{\partial_{11} v_2^{k-2s}}{2s+3} - \sum_{s=0}^{\lfloor \frac{k-1}{2} \rfloor} \frac{\partial_{11} u_2^{k-2s-1}}{2s+3} \right), \\
 X_6^k &= X_6 + \lambda \left(v_3^{k-1} + \partial_2 u_2^{k-1} \right) - \frac{\mu}{3} \partial_{22} v_3^{k-1} + \\
 &\quad + \sum_{s=0}^{\lfloor \frac{k-1}{2} \rfloor} \left(\frac{\mu}{2s+3} \partial_{11} v_3^{k-2s-1} + \frac{\lambda + 3\mu}{2s+1} \partial_1 u_1^{k-2s-1} \right) - (\lambda + 2\mu) \sum_{s=1}^{\lfloor \frac{k-1}{2} \rfloor} \frac{v_3^{k-2s-1}}{2s+1} - \\
 &\quad - \sum_{s=1}^{\lfloor \frac{k}{2} \rfloor} \left(\frac{\mu}{2s+1} \partial_{11} u_3^{k-2s} + \frac{\lambda + 3\mu}{2s+1} \partial_1 v_1^{k-2s} - \frac{\lambda + 2\mu}{2s-1} u_3^{k-2s} \right), \\
 &\quad \left(k = 1, 2, \dots; \quad u_3^k = v_3^k = u_\alpha^k = v_\alpha^k = 0, \text{ if } k < 0; \quad \alpha = 1, 2, \right. \\
 &\quad \left. \theta = \partial_1 u_1^k + \partial_2 u_2^k, \quad \Theta = \partial_1 v_1^k + \partial_2 v_2^k, \quad \partial_\alpha = \frac{\partial}{\partial x^\alpha} \right).
 \end{aligned}$$

Here λ and μ are Lamé's constants, X_i ($i = 1, \dots, 6$) - the components of external force.

For each fixed k equations (1) and (2) coincide with equations of I. N. Vekua which is obtained for prismatic shells with constant thickness in case of approximation $N = 1$. The right part of equations (1) and (2) is well-known quantity, defined by functions $u_i^0, \dots, u_i^{k-1}, v_j^0, \dots, v_j^{k-1}$ ($i, j = 1, 2, 3$).

The general solutions of systems (1) and (2) are written in the following form:

$$\begin{aligned} 2\mu u_+^k &= \varkappa \varphi^k(z) - z \overline{\varphi'^k(z)} - \overline{\psi^k(z)} - \frac{\lambda}{6(\lambda + \mu)} \frac{\partial \chi^k(z, \bar{z})}{\partial \bar{z}} + \widehat{u}_+^k, \\ 2\mu v_3^k &= -\frac{2\lambda}{3\lambda + 2\mu} \left(\overline{\varphi'^k(z)} + \overline{\varphi'^k(z)} \right) + \chi^k(z, \bar{z}) + \widehat{v}_3^k, \\ 2\mu v_+^k &= \frac{4(\lambda + 2\mu)}{3\mu} \overline{f''^k(z)} + z \overline{f'^k(z)} + \overline{f^k(z)} - \overline{2g'^k(z)} + i \frac{\partial w^k(z, \bar{z})}{\partial \bar{z}} + \widehat{v}_+^k, \\ 2\mu u_3^k &= -\frac{1}{2} \left(\bar{z} \overline{f^k(z)} + \overline{z f^k(z)} \right) + \overline{g^k(z)} + \overline{g^k(z)} + \widehat{u}_3^k, \\ &\left(u_+^k = u_1^k + i u_2^k, \quad v_+^k = v_1^k + i v_2^k, \quad z = x^1 + i x^2, \right. \\ &\left. \frac{\partial}{\partial \bar{z}} = \frac{1}{2} \left(\frac{\partial}{\partial x^1} + i \frac{\partial}{\partial x^2} \right), \quad \frac{\partial}{\partial z} = \frac{1}{2} \left(\frac{\partial}{\partial x^1} - i \frac{\partial}{\partial x^2} \right) \right), \end{aligned}$$

where $\varkappa = \frac{5\lambda + 6\mu}{3\lambda + 2\mu}$, $\varphi^k(z)$, $\psi^k(z)$, $f^k(z)$ and $g^k(z)$ are any analytic functions of z , $\chi^k(z, \bar{z})$ and $w^k(z, \bar{z})$ are the general solutions of the following Helmholtz's equations, respectively:

$$\begin{aligned} \Delta \chi^k - \eta^2 \chi^k &= 0 \quad \left(\eta^2 = \frac{12(\lambda + \mu)}{\lambda + 2\mu} \right), \\ \Delta w^k - \gamma^2 w^k &= 0 \quad (\gamma^2 = 3). \end{aligned}$$

Here \widehat{u}_+^k , \widehat{v}_3^k and \widehat{v}_+^k , \widehat{u}_3^k are particular solutions of the non-homogeneous equations (1) and (2), respectively.

Let's consider the following boundary value problems for system (1),(2).

Find the solutions of the homogeneous system of equations (1) and (2) compatible with the kinematic boundary conditions:

$$u_+^k = \begin{cases} F_+^k, & |z| = R_1, \\ F_+^k, & |z| = R_2, \end{cases} \quad v_3^k = \begin{cases} F_3^k, & |z| = R_1, \\ F_3^k, & |z| = R_2, \end{cases} \quad (3)$$

$$v_+^k = \begin{cases} G_+^k, & |z| = R_1, \\ G_+^k, & |z| = R_2, \end{cases} \quad u_3^k = \begin{cases} G_3^k, & |z| = R_1, \\ G_3^k, & |z| = R_2, \end{cases} \quad (4)$$

where F_+, F_3, G_+, G_3 are the known functions of $u_i, \dots, u_i^{k-1}, v_j, \dots, v_j^{k-1}$ ($i, j, = 1, 2, 3$).

Let us introduce the functions $\varphi(z), \psi(z), \chi(z, \bar{z}), F_+$ and F_3 by the series

$$\begin{aligned} \varphi(z) &= Az \ln z + a \ln z + \sum_{n=1}^{\infty} a_n z^n, & \psi(z) &= \beta \ln z + \sum_{n=0}^{\infty} b_n z^n, \\ \chi(z, \bar{z}) &= \sum_{-\infty}^{\infty} \left(\alpha_n I_n(\eta r) + \alpha_n' K_n(\eta r) \right) e^{in\theta}, \\ F_+ &= \sum_{-\infty}^{\infty} A_n e^{in\theta}, & F_3 &= \sum_{-\infty}^{\infty} B_n e^{in\theta}, \end{aligned} \quad (5)$$

where $I_n(\eta r)$ and $K_n(\eta r)$ are Bessel's modified functions.

By substituting (5) into (3) we obtain the system of algebraic equations:

$$\begin{aligned} A &= 0, & \alpha &+ \bar{\beta} = 0, \\ 2\alpha \ln R_1 \alpha - 2R_1^2 \bar{a}_2 - \frac{\lambda \eta}{12(\lambda + \mu)} \left(I_0(\eta R_1) \alpha_{-1}^k - K_0(\eta R_1) \alpha_{-1}^{k'} \right) + \alpha a_0 - b_0 &= A', \\ 2\alpha \ln R_2 \alpha - 2R_2^2 \bar{a}_2 - \frac{\lambda \eta}{12(\lambda + \mu)} \left(I_0(\eta R_2) \alpha_{-1}^k - K_0(\eta R_2) \alpha_{-1}^{k'} \right) + \alpha a_0 - b_0 &= A'', \\ -\bar{\alpha} + \alpha R_1^2 a_2 - \frac{\lambda \eta}{12(\lambda + \mu)} \left(I_2(\eta R_1) \alpha_1^k - K_2(\eta R_1) \alpha_1^{k'} \right) - R_1^{-2} \bar{b}_{-2} &= A', \\ -\bar{\alpha} + \alpha R_2^2 a_2 - \frac{\lambda \eta}{12(\lambda + \mu)} \left(I_2(\eta R_2) \alpha_1^k - K_2(\eta R_2) \alpha_1^{k'} \right) - R_2^{-2} \bar{b}_{-2} &= A'', \\ \alpha R_1^n a_n + (n-2) R_1^{-n+2} \bar{a}_{-n+2} - \frac{\lambda \eta}{12(\lambda + \mu)} \left(I_n(\eta R_1) \alpha_{n-1}^k - K_n(\eta R_1) \alpha_{n-1}^{k'} \right) & \\ & - R_1^{-n} \bar{b}_{-n} = A', \\ \alpha R_2^n a_n + (n-2) R_2^{-n+2} \bar{a}_{-n+2} - \frac{\lambda \eta}{12(\lambda + \mu)} \left(I_n(\eta R_2) \alpha_{n-1}^k - K_n(\eta R_2) \alpha_{n-1}^{k'} \right) & \\ & - R_2^{-n} \bar{b}_{-n} = A'', \\ I_1(\eta R_1) \alpha_1^k + K_1(\eta R_1) \alpha_1^{k'} - \frac{2\lambda}{3\lambda + 2\mu} \left(2R_1 a_2 - \frac{\bar{\alpha}}{R_1} \right) &= B', \\ I_1(\eta R_2) \alpha_1^k + K_1(\eta R_2) \alpha_1^{k'} - \frac{2\lambda}{3\lambda + 2\mu} \left(2R_2 a_2 - \frac{\bar{\alpha}}{R_2} \right) &= B'', \\ I_n(\eta R_1) \alpha_n^k + K_n(\eta R_1) \alpha_n^{k'} - \frac{2\lambda}{3\lambda + 2\mu} \left((n+1) R_1^n a_{n+1} - (n-1) R_1^{-n} \bar{a}_{-n+1} \right) &= B'_n, \\ I_n(\eta R_2) \alpha_n^k + K_n(\eta R_2) \alpha_n^{k'} - \frac{2\lambda}{3\lambda + 2\mu} \left((n+1) R_2^n a_{n+1} - (n-1) R_2^{-n} \bar{a}_{-n+1} \right) &= B''_n, \\ (n = \pm 1, -2, \pm 3, \dots). & \end{aligned} \quad (6)$$

For coefficients a_n and α we have:

$$\begin{aligned} a_2 &= \frac{E_1 C_2 - E_2 C_0}{E_1 T_2 - E_2 S_0}, & \bar{\alpha} &= \frac{C_0 T_2 - C_2 S_0}{E_1 T_2 - E_2 S_0}, \\ a_n &= \frac{T_{-n+2} C_n - S_n C_{-n+2}}{T_n T_{-n+2} - S_n S_{-n+2}}, \end{aligned}$$

where

$$\begin{aligned}
E_1 &= 2\alpha \ln \frac{R_2}{R_1} - \frac{\lambda^2 \eta (I_0(\eta R_2) - I_0(\eta R_1))(R_1 K_1(\eta R_1) - R_2 K_1(\eta R_2))}{6(\lambda + \mu)(3\lambda + 2\mu)R_1 R_2 (I_1(\eta R_2) K_1(\eta R_1) - I_1(\eta R_1) K_1(\eta R_2))} \\
&\quad + \frac{\lambda^2 \eta (K_0(\eta R_2) - K_0(\eta R_1))(R_1 I_1(\eta R_1) - R_2 I_1(R_2))}{6(\lambda + \mu)(3\lambda + 2\mu)R_1 R_2 (I_1(\eta R_2) K_1(\eta R_1) - I_1(\eta R_1) K_1(\eta R_2))}, \\
E_2 &= R_1^2 - R_2^2 - \frac{\lambda^2 \eta (R_2^2 I_2(\eta R_2) - R_1^2 I_2(\eta R_1))(R_1 K_1(\eta R_1) - R_2 K_1(\eta R_2))}{6(\lambda + \mu)(3\lambda + 2\mu)R_1 R_2 (I_1(\eta R_2) K_1(\eta R_1) - I_1(\eta R_1) K_2(\eta R_2))} \\
&\quad + \frac{\lambda^2 \eta (R_2^2 K_2(\eta R_2) - R_1^2 K_2(\eta R_1))(R_1 I_1(\eta R_1) - R_2 I_1(R_2))}{6(\lambda + \mu)(3\lambda + 2\mu)R_1 R_2 (I_1(\eta R_2) K_1(\eta R_1) - I_1(\eta R_1) K_1(\eta R_2))}, \\
C_n &= A_n'' R_2^n - A_n' R_1^n \\
&\quad + \frac{\lambda \eta (R_2^n I_n(\eta R_2) - R_1^n I_n(\eta R_1))(B_{n-1}'' K_{n-1}(\eta R_1) - B_{n-1}' K_{n-1}(\eta R_2))}{12(\lambda + \mu)(I_{n-1}(\eta R_2) K_{n-1}(\eta R_1) - I_{n-1}(\eta R_1) K_{n-1}(\eta R_2))} \\
&\quad - \frac{\lambda \eta (R_2^n K_n(\eta R_2) - R_1^n K_n(\eta R_1))(B_{n-1}'' I_{n-1}(\eta R_1) - B_{n-1}' I_{n-1}(\eta R_2))}{12(\lambda + \mu)(I_{n-1}(\eta R_2) K_{n-1}(\eta R_1) - I_{n-1}(\eta R_1) K_{n-1}(\eta R_2))}, \\
T_n &= \alpha (R_2^{2n} - R_1^{2n}) \\
&\quad - \frac{\lambda^2 \eta n (R_2^n I_n(\eta R_2) - R_1^n I_n(\eta R_1))(R_2^{n-1} K_{n-1}(\eta R_1) - R_1^{n-1} K_{n-1}(\eta R_2))}{6(\lambda + \mu)(3\lambda + 2\mu)(I_{n-1}(\eta R_2) K_{n-1}(\eta R_1) - I_{n-1}(\eta R_1) K_{n-1}(\eta R_2))} \\
&\quad + \frac{\lambda^2 \eta n (R_2^n K_n(\eta R_2) - R_1^n K_n(\eta R_1))(R_2^{n-1} I_{n-1}(\eta R_1) - R_1^{n-1} I_{n-1}(\eta R_2))}{6(\lambda + \mu)(3\lambda + 2\mu)(I_{n-1}(\eta R_2) K_{n-1}(\eta R_1) - I_{n-1}(\eta R_1) K_{n-1}(\eta R_2))}, \\
S_n &= (n-2) [R_2^2 - R_1^2 \\
&\quad + \frac{\lambda^2 \eta (R_2^n I_n(\eta R_2) - R_1^n I_n(\eta R_1))(R_2^{-n+1} K_{n-1}(\eta R_1) - R_1^{-n+1} K_{n-1}(\eta R_2))}{6(\lambda + \mu)(3\lambda + 2\mu)(I_{n-1}(\eta R_2) K_{n-1}(\eta R_1) - I_{n-1}(\eta R_1) K_{n-1}(\eta R_2))} \\
&\quad - \frac{\lambda^2 \eta (R_2^n K_n(\eta R_2) - R_1^n K_n(\eta R_1))(R_2^{-n+1} I_{n-1}(\eta R_1) - R_1^{-n+1} I_{n-1}(\eta R_2))}{6(\lambda + \mu)(3\lambda + 2\mu)(I_{n-1}(\eta R_2) K_{n-1}(\eta R_1) - I_{n-1}(\eta R_1) K_{n-1}(\eta R_2))}].
\end{aligned}$$

Let

$$\begin{aligned}
f(z) &= Cz \ln z + \gamma \ln z + \sum_{n=1}^{\infty} c_n z^n, \quad g(z) = \delta \ln z + \sum_{n=0}^{\infty} d_n z^n, \\
w(z, \bar{z}) &= \sum_{-\infty}^{\infty} \left(\beta_n I_n(\gamma r) + \beta_n' K_n(\gamma r) \right) e^{in\theta}, \\
G_+ &= \sum_{-\infty}^{\infty} M_n e^{in\theta}, \quad G_- = \sum_{-\infty}^{\infty} N_n e^{in\theta}.
\end{aligned}$$

We now find the coefficients c_n , d_n , β_n and β_n' from following system of algebraic equations

$$\begin{aligned}
\frac{i\gamma}{2} \left(I_1(\gamma R_1) \beta_0^k - K_1(\gamma R_1) \beta_0'^k \right) + R_1 \left(c_1^k + \bar{c}_1^k \right) - \frac{2\bar{\delta}}{R_1} &= M_1', \\
\frac{i\gamma}{2} \left(I_1(\gamma R_2) \beta_0^k - K_1(\gamma R_2) \beta_0'^k \right) + R_2 \left(c_1^k + \bar{c}_1^k \right) - \frac{2\bar{\delta}}{R_2} &= M_1'',
\end{aligned}$$

$$\begin{aligned} & \frac{i\gamma}{2} \left(I_n(\gamma R_1) \beta_{n-1}^k - K_1(\gamma R_1) \beta_{n-1}^{k'} \right) + R_1^n c_n^k + \frac{4(\lambda + 2\mu)}{3\mu} (n-1)(n-2) R_1^{-n} \bar{c}_{-n+2}^{\bar{k}} \\ & - (n-2) R_1^{-n+2} \bar{c}_{-n+2}^{\bar{k}} + 2(n-1) R_1^{-n} \bar{d}_{-n+1} = M'_n, \\ & \frac{i\gamma}{2} \left(I_n(\gamma R_2) \beta_{n-1}^k - K_1(\gamma R_2) \beta_{n-1}^{k'} \right) + R_2^n c_n^k + \frac{4(\lambda + 2\mu)}{3\mu} (n-1)(n-2) R_2^{-n} \bar{c}_{-n+2}^{\bar{k}} \\ & - (n-2) R_2^{-n+2} \bar{c}_{-n+2}^{\bar{k}} + 2(n-1) R_2^{-n} \bar{d}_{-n+1} = M''_n, \\ & d_0^k + \bar{d}_0^{\bar{k}} - \frac{R_1^2}{2} \left(c_1^k + \bar{c}_1^{\bar{k}} \right) + 2\delta \ln R_1 = N'_0, \\ & d_0^k + \bar{d}_0^{\bar{k}} - \frac{R_2^2}{2} \left(c_1^k + \bar{c}_1^{\bar{k}} \right) + 2\delta \ln R_2 = N''_0, \\ & R_1^n d_n^k + R_1^{-n} \bar{d}_{-n}^{\bar{k}} - \frac{1}{2} \left(R_1^{n+2} c_{n+1}^k + R_1^{-n+2} \bar{c}_{-n+1}^{\bar{k}} \right) = N'_n, \\ & R_2^n d_n^k + R_2^{-n} \bar{d}_{-n}^{\bar{k}} - \frac{1}{2} \left(R_2^{n+2} c_{n+1}^k + R_2^{-n+2} \bar{c}_{-n+1}^{\bar{k}} \right) = N''_n. \end{aligned}$$

R E F E R E N C E S

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