Reports of Enlarged Session of the Seminar of I. Vekua Institute of Applied Mathematics Volume 20, $N^{0}2$, 2005

TWO ELASTIC SOLIDS INTERACTION PROBLEM

Khmiadashvili M., Skhvitaridze K.

Georgian Technical University

Abstract

The contact problem for homogeneous equation systems of classic and asymmetrical theory is considered, when a spare is the contact surface. it is proved, that the problem has a singular solution. The solution is presented in the form of absolutely and uniformly convergent series.

Keywords and phrases: continual theory, asymmetrical theory, Green's formula, metaharmonic function.

AMS subject classification (2000): 35J55; 75H20; 5H25.

Let Ω^+ be the spere bounded by spherical surface $\partial \Omega$ with the center in the origin and the radius R.

Problem (C). Find the regular vectors $u^{(0)}(x)$ and $u^{(1)}(x)$ in Ω^+ and Ω^- areas accordingly, satisfying:

1. equations [1.3]

$$M(\partial x)u^{(o)}(x) = \Delta(\mu_0 - \nu_0 \Delta)u^{(0)}(x) + (\lambda_0 + \mu_0 + \nu_0 \Delta) \text{ grad div } u^{(0)}(x) = 0, \quad x \in \Omega^+,$$

$$A(\partial x)u^{(1)}(x) = \mu_1 \Delta u^{(1)}(x) + (\lambda_1 + \mu_1) \text{ grad div } u^{(1)}(x) = 0, \quad x \in \Omega^-,$$
(1)

2. boundary conditions

$$[u^{(0)}(z)]^{+} - [u^{(1)}(z)]^{-} = f^{(1)}(z), \quad z \in \partial\Omega,$$

$$[T^{(0)}(\partial z, n)u^{(0)}(z)]^{+} - [T^{(1)}(\partial z, n)u^{(1)}(z)]^{-} = f^{(2)}(z), \quad z \in \partial\Omega,$$

$$[R^{(0)}(\partial z, n)u^{(0)}(z)]^{+} = f^{(3)}(z), \quad z \in \partial\Omega,$$
(2)

3. Vector $u^{(1)}(x)$ satisfies nearby infinitely distant pint the conditions:

$$u_j^{(1)}(x) = O(|x|^{-1}), \quad \frac{\partial u_j^{(1)}(x)}{\partial x_k} = o(|x|^{-1}), \quad k, j = 1, 2, 3,$$
 (3)

where $f^{(j)} = (f_1^{(j)}, f_2^{(j)}, f_3^{(j)},), \quad j = 1, 2, 3$ are the given vectors on $\partial\Omega$, n(z) is the ort of normal outside with regard to Ω^+ in point $z \in \partial\Omega$. $u^{(j)}(x) = (u_1^{(j)}(x), u_2^{(j)}(x), u_3^{(j)}(x)), \quad j = 0, 1$ - the vector if displacement, Δ -laplace operator.

$$T^{(0)}(\partial x, n)u^{(0)} = 2\mu_0 \frac{\partial u^{(0)}}{\partial n} + \lambda_0 n \operatorname{div} u^{(0)} + \mu_0 [n \times \operatorname{rot} u^{(0)}] + \nu_0 [n \times \Delta \operatorname{rot} u^{(0)}] - \frac{1}{2} [n \times \operatorname{grad}(n \cdot q(u^{(0)}))],$$

$$T^{(0)}(\partial x, n)u^{(1)} = 2\mu_1 \frac{\partial u^{(1)}}{\partial n} + \lambda_1 n \operatorname{div} u^{(1)} + \mu_1 [n \times \operatorname{rot} u^{(1)}],$$

$$R(\partial x, n)u^{(0)} = \operatorname{rot} u^{(0)} - n(n \cdot \operatorname{rot} u^{(0)}),$$

$$q(u^{(0)}) = 2(\nu_0 + \nu'_0)\frac{\partial}{\partial n} \operatorname{rot} u^{(0)} + 2\nu' [n \times \operatorname{rot} \operatorname{rot} u^{(0)}],$$

 $\lambda_j, \ \mu_j, \ \nu_0, \ \nu_0', \ \ j = 0, 1$ -elastic constants, satisfying the following conditions:

 $\mu_j > 0, \quad 3\lambda_j + 2\mu_j > 0, \quad |\nu'_0| \le \nu_0, \quad j = 0, 1.$

Theorem 1. The problem (C) permits not more than one regular solution.

Proof. The theorem will be proved, if we show that the corresponding homogeneous problem $(f^{(j)}(z) = 0, j = 1, 2, 3)$ has only trivial solution. Green's formula in Ω^- domain area for system (1) has a form [3,5]

$$\int_{\Omega^{-}} u^{(1)}(x) \cdot A(\partial x) u^{(1)}(x) \, dx = -\int_{\partial\Omega} [u^{(1)}(z)]^{-} \cdot [T^{(1)}(\partial z, n) u^{(1)}(z)]^{-} ds \qquad (4)$$
$$-\int_{\Omega^{-}} \widetilde{E}^{(1)}(u^{(1)}, u^{(1)}) \, dx,$$

keeping in mind boundary conditions of problem $(C)_0$, we'll have:

$$\int_{\Omega^{+}} u^{(0)}(x) \cdot M(\partial x) u^{(0)}(x) \, dx = \int_{\partial \Omega} \{ [u^{(0)}(z)]^{+} \cdot [T^{(0)}(\partial z, n) u^{(0)}(z)]^{+} + [R(\partial z, n) u^{(0)}(z)]^{+} \cdot [Q^{(0)}(\partial z, n) u^{(0)}(z)]^{+} \} \, ds - \int_{\Omega^{+}} E^{(0)}(u^{(0)}, u^{(0)}) \, dx,$$
(5)

where [3,5]

$$\begin{aligned} \widetilde{E}^{(1)}(u^{(1)}, u^{(1)}) &\geq 0, \quad E^{(0)}(u^{(0)}, u^{(0)}) \geq 0, \\ \omega_k^{(0)} &= \frac{1}{2} (\operatorname{rot} u^{(0)})_k, \quad k = 1, 2, 3, \\ Q^{(0)}(\partial x, n) u^{(0)} &= \frac{1}{2} \Big[q(u^{(0)}) - n(n \cdot q(u^{(0)})) \Big]. \end{aligned}$$

Keeping in mind boundary conditions of problem $(C)_0$, we'll have:

$$E^{(0)}(u^{(0)}, u^{(0)}) = 0, \quad \widetilde{E}^{(1)}(u^{(1)}, u^{(1)}) = 0.$$

The solutions of these equations have the form [3,5]:

$$\begin{split} u^{(0)}(x) &= [a^{(0)} \times x] + b^{(0)}, \quad x \in \Omega^+, \\ u^{(1)}(x) &= [a^{(1)} \times x] + b^{(1)}, \quad x \in \Omega^-, \end{split}$$

Taking into consideration conditions (3) and boundary conditions of problem $(C)_0$, we'll have: $u^{(0)}(x) = 0$, $x \in \Omega^+$, $u^{(0)}(x) = 0$, $x \in \Omega^-$. \Box

The solution of problem (C) we'll search in the form [5]:

$$u^{(0)}(x) = \operatorname{grad} \Phi_1^{(0)}(x) - a_0 \operatorname{grad} r^2 \left(r \frac{\partial}{\partial r} + 1 \right) \Phi_2^{(0)}(x) + \operatorname{rot} \operatorname{rot} (xr^2 \Phi_2^{(0)}(x)) + \operatorname{rot} (x\Phi_3^{(0)}(x)) + \operatorname{rot} \operatorname{rot} (x\Phi_4^{(0)}(x)) + \operatorname{rot} (x\Phi_5^{(0)}(x)), \quad x \in \Omega^+,$$
(6)
$$u^{(1)}(x) = \operatorname{grad} \Phi_1^{(1)}(x) - a_1 \operatorname{grad} r^2 \left(r \frac{\partial}{\partial r} + 1 \right) \Phi_2^{(1)}(x) + \operatorname{rot} \operatorname{rot} (xr^2 \Phi_2^{(1)}(x)) + \operatorname{rot} (x\Phi_3^{(1)}(x)) \quad x \in \Omega^+,$$

where $\Phi_j^{(l)}(x)$, l = 1, 2, j = 1, 2, 3-scalar harmonic functions, $\Phi_j^{(0)}(x)$, j = 4, 5-scalar metaharmonic functions.

$$(\Delta - l^2) \Phi_j^{(0)}(x) = 0, \quad j = 4, 5, \quad l^2 = \frac{\mu}{\nu}, a_l = \mu (\lambda_l + 2\mu_l)^{-1}, \quad l = 0, 1,$$

 $x = (x_1, x_2, x_3), \quad r = |x|, \quad r \frac{\partial}{\partial r} = x \cdot \text{grad}.$

We write down functions $\Phi_j^{(l)}(x)$ in the form:

$$\Phi_{j}^{(0)}(x) = \sum_{k=0}^{\infty} \sum_{m=-k}^{k} \left(\frac{r}{R}\right)^{k} Y_{k}^{(m)}(\theta,\varphi) A_{mk}^{(j)}, \quad j = 1, 2, 3,$$

$$\Phi_{j}^{(0)}(x) = \sum_{k=0}^{\infty} \sum_{m=-k}^{k} g_{k}(lr) Y_{k}^{(m)}(\theta,\varphi) A_{mk}^{(j)}, \quad j = 4, 5,$$

$$\Phi_{j}^{(1)}(x) = \sum_{k=0}^{\infty} \sum_{m=-k}^{k} \left(\frac{R}{r}\right)^{k} Y_{k}^{(m)}(\theta,\varphi) B_{mk}^{(j)}, \quad j = 1, 2, 3,$$
(7)

where (r, θ, φ) -spherical coordinates of the point x; $A_{mk}^{(j)}, B_{mk}^{(j)}$ -desired constants

$$Y_k^{(m)}(\theta,\varphi) = \sqrt{\frac{2k+1}{4\pi} \cdot \frac{(k-m)!}{(k+m)!}} P_k^{(m)}(\cos\theta) e^{im\varphi},$$
$$g_k(lr) = \sqrt{\frac{R}{r}} \frac{I_{k+\frac{1}{2}}(lr)}{I_{k+\frac{1}{2}}(lR)},$$

 $P_k^{(m)}(\cos\theta)\text{-Lejandr's}$ added function of the first type, $I_{k+\frac{1}{2}}(lr)\text{-Besel's}$ function of imaginary argument.

if the value of functions $\Phi_j^{(l)}(x)$ put from (7) to (6), we'll have:

$$u^{(j)}(x) = \sum_{k=0}^{\infty} \sum_{m=-k}^{k} \{u_{mk}^{(j)}(r) X_{mk}(\theta,\varphi) + \sqrt{k(k+1)} [v_{mk}^{(j)}(r) Y_{mk}(\theta,\varphi) + \sqrt{k(k+1$$

Where $u_{mk}^{(j)}(r), \ldots, c_{mk}^{(j)}(r), \quad b_{mk}^{(2)}(r), \quad c_{mk}^{(2)}(r), \quad j = 0, 1 \text{-functions of } r, \quad X_{mk}(\theta, \varphi),$ $Y_{mk}(\theta, \varphi), \quad Z_{mk}(\theta, \varphi) \text{-ortonormalized vectors in } L_2(\Sigma_1) \text{ class } [4].$

Let vector $f^{(j)}(z)$, j = 1, 2, 3 satisfy the conditions, under which we can spread it out by Fourie-Laplass series in system

$$\{X_{mk}(\theta,\varphi), Y_{mk}(\theta,\varphi), Z_{mk}(\theta,\varphi)\}_{|m| \le k, k=\overline{0,\infty},}$$

$$f^{(j)}(z) = \sum_{k=o}^{\infty} \sum_{m=-k}^{k} \{\alpha_{mk}^{(j)} X_{mk}(\theta,\varphi) + \sqrt{k(k+1)} [\beta_{mk}^{(j)} Y_{mk}(\theta,\varphi) + \gamma_{mk}^{(j)} Z_{mk}(\theta,\varphi)]\}(9)$$

$$f^{(3)}(z) = \sum_{k=1}^{\infty} \sum_{m=-k}^{k} \sqrt{k(k+1)} [\beta_{mk}^{(3)} Y_{mk}(\theta,\varphi) + \gamma_{mk}^{(3)} Z_{mk}(\theta,\varphi)]. \quad j = 1, 2,$$

proceeding to limit in both sides of the equality (8), when $x \to z \in \partial \Omega$ keeping in mind boundary conditions (2), formulas (9) also, we'll have following system for desired constants:

$$\begin{split} u_{mk}^{(0)}(R) &- u_{mk}^{(1)}(R) = \alpha_{mk}^{(1)}, \quad a_{mk}^{(0)}(R) - a_{mk}^{(1)}(R) = \alpha_{mk}^{(2)}, \quad k \ge 0, \\ v_{mk}^{(0)}(R) &- v_{mk}^{(1)}(R) = \beta_{mk}^{(1)}, \quad b_{mk}^{(0)}(R) - b_{mk}^{(1)}(R) = \beta_{mk}^{(2)}, \\ w_{mk}^{(0)}(R) &- w_{mk}^{(1)}(R) = \gamma_{mk}^{(1)}, \quad c_{mk}^{(0)}(R) - c_{mk}^{(1)}(R) = \gamma_{mk}^{(2)}, \\ b_{mk}^{(2)}(R) &= \beta_{mk}^{(3)}, \quad c_{mk}^{(2)}(R) = \gamma_{mk}^{(3)}, \quad k \ge 1. \end{split}$$

This system is compatible according to uniqueness theorem, Putting the solution of the system (9) in (8), we'll find the solution of the problem (C). if $f^{(1)}(z) \in C^4(\partial\Omega)$, $f^{(j)}(z) \in C^3(\partial\Omega)$, j = 2, 3, then series (8) are absolutely and uniformly convergent in $\overline{\Omega^+}$ and $\overline{\Omega^-}$ areas.

REFERENCES

1. Cosserat E., Cosserat F. Therie des corps deformables, Herman, Paris, 1909.

2. Giorgashvili l. It is proved, that the problem has a singular solution. it is proved, that the problem has a singular solution, ., Solution of the basic boundary value problems of stationary

thermoelastic oscillations for domains bounded by spherical surfaces, Georgian Math. J. 4 (1997), N 5, p. 421–438.

3. Kupradze V.D., Gegelia T.G., Basheleishvili M.O., Burchuladze T.V. Three-dimensional problems of the mathematical theory of elasticity and thermoelasticity, Nauka, Moscow, 1976.

4. Mors F. M., Feshbakh G. Methods of Theoretical Physics, Foreign Literature Publishing House. Moscow. 2. 1980.

5. Khmiadashvili M. The solution of basic static boundary value problems in asymmetrical theory of elasticity for the spere, Inst. Appl. mathem. 46 (1992), 215-226.

Received 26. IV. 2005; accepted 20. IX. 2005.