

ASYMPTOTIC BEHAVIOR OF THE SOLUTION AND  
SEMI-DISCRETE FINITE DIFFERENCE SCHEME FOR ONE  
NONLINEAR INTEGRO-DIFFERENTIAL MODEL WITH SOURCE  
TERMS

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*Abstract*

One nonlinear integro-differential system with source terms is considered. The model arises at describing penetration of a magnetic field into a substance. Large time behavior of solution of the initial-boundary value problem is given. Corresponding semi-discrete finite difference scheme is studied as well.

*Key words and phrases:* Nonlinear integro-differential system, asymptotic behavior, semi-discrete scheme.

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## 1 Introduction

One system of nonlinear integro-differential equations is considered. Large time behavior of solution and semi-discrete finite difference scheme for the initial-boundary value problem is studied. The investigated system arises in mathematical modeling of the process of a magnetic field penetration into a substance. If the coefficient of thermal heat capacity and electroconductivity of the substance highly dependent on temperature, then the Maxwell's system [1], that describes above-mentioned process, can be rewritten in the following form [2]:

$$\frac{\partial H}{\partial t} = -rot \left[ a \left( \int_0^t |rot H|^2 d\tau \right) rot H \right], \quad (1.1)$$

where  $H = (H_1, H_2, H_3)$  is a vector of the magnetic field and the function  $a = a(S)$  is defined for  $S \in [0, \infty)$ .

Note that (1.1) is complex. Special cases of such type models were investigated (see, for example, [2]-[12] and references therein). Investigations mainly are carried out for one-component magnetic field cases. The existence of global solutions for initial-boundary value problems of such models have been proven in [2]-[5],[11] by using the Galerkin and compactness methods [13],[14]. The asymptotic behavior of the solutions have been the subject of intensive research as well (see, for example, [11],[15],[16] and references therein).

The following one-dimensional system with two-component magnetic field is considered in many works as well (see, for example, [17]-[22]):

$$\begin{aligned} \frac{\partial U}{\partial t} &= \frac{\partial}{\partial x} \left\{ a \left( \int_0^t \left[ \left( \frac{\partial U}{\partial x} \right)^2 + \left( \frac{\partial V}{\partial x} \right)^2 \right] d\tau \right) \frac{\partial U}{\partial x} \right\}, \\ \frac{\partial V}{\partial t} &= \frac{\partial}{\partial x} \left\{ a \left( \int_0^t \left[ \left( \frac{\partial U}{\partial x} \right)^2 + \left( \frac{\partial V}{\partial x} \right)^2 \right] d\tau \right) \frac{\partial V}{\partial x} \right\}, \end{aligned} \tag{1.2}$$

where  $a = a(S)$  is a given function.

For the system (1.2) the convergence of the semi-discrete and full finite difference approximations of the initial-boundary value problem for the case  $a(S) = 1 + S$  with first kind boundary conditions were studied in [22].

The aim of this note is to study asymptotic behavior of solution as  $t \rightarrow \infty$  and to construct semi-discrete approximate solutions for one generalization of the system type (1.2) by adding monotonic nonlinear source terms.

## 2 Statement of Problem and Main Results

In the  $[0, 1] \times [0, \infty)$  let us consider following initial-boundary value problem:

$$\begin{aligned} \frac{\partial U}{\partial t} &= \frac{\partial}{\partial x} \left\{ \left( 1 + \int_0^t \left[ \left( \frac{\partial U}{\partial x} \right)^2 + \left( \frac{\partial V}{\partial x} \right)^2 \right] d\tau \right)^p \frac{\partial U}{\partial x} \right\} - |U|^{q-2} U, \\ \frac{\partial V}{\partial t} &= \frac{\partial}{\partial x} \left\{ \left( 1 + \int_0^t \left[ \left( \frac{\partial U}{\partial x} \right)^2 + \left( \frac{\partial V}{\partial x} \right)^2 \right] d\tau \right)^p \frac{\partial V}{\partial x} \right\} - |V|^{q-2} V, \end{aligned} \tag{2.1}$$

$$U(0, t) = U(1, t) = V(0, t) = V(1, t) = 0, \tag{2.2}$$

$$U(x, 0) = U_0(x), \quad V(x, 0) = V_0(x), \quad (2.3)$$

where  $0 < p \leq 1$ ,  $q \geq 2$ ;  $U_0 = U_0(x)$  and  $V_0 = V_0(x)$  are given functions.

The following statement is true.

**Theorem 1.** *If  $0 < p \leq 1$ ,  $q \geq 2$  and  $U_0, V_0 \in H_0^1(0, 1)$ , then problem (2.1) - (2.3) has not more than one solution and the following asymptotic property takes place*

$$\|U\| + \left\| \frac{\partial U}{\partial x} \right\| + \|V\| + \left\| \frac{\partial V}{\partial x} \right\| \leq C \exp\left(-\frac{t}{2}\right).$$

Here  $\|\cdot\|$  is the usual norm of the space  $L_2(0, 1)$  and  $C$  denotes positive constant independent of  $t$ .

On  $[0, 1]$  let us introduce a net with mesh points denoted by  $x_i = ih$ ,  $i = 0, 1, \dots, M$ , with  $h = 1/M$ . The boundaries are specified by  $i = 0$  and  $i = M$ . The semi-discrete approximation at  $(x_i, t)$  is designed by  $u_i = u_i(t)$  and  $v_i = v_i(t)$ . The exact solution to the problem at  $(x_i, t)$  is denoted by  $U_i = U_i(t)$  and  $V_i = V_i(t)$ . At points  $i = 1, 2, \dots, M - 1$ , the integro-differential equation will be replaced by approximation of the space derivatives by a forward and backward differences. We will use the following known notations:

$$r_{x,i}(t) = \frac{r_{i+1}(t) - r_i(t)}{h}, \quad r_{\bar{x},i}(t) = \frac{r_i(t) - r_{i-1}(t)}{h}.$$

Using usual methods of construction of discrete analogs (see, for example, [26]) let us construct the following semi-discrete finite difference scheme for problem (2.1) - (2.3):

$$\begin{aligned} \frac{du_i}{dt} &= \left\{ \left( 1 + \int_0^t [(u_{\bar{x},i})^2 + (v_{\bar{x},i})^2] d\tau \right)^p u_{\bar{x},i} \right\}_x - |u_i|^{q-2} u_i, \\ \frac{dv_i}{dt} &= \left\{ \left( 1 + \int_0^t [(u_{\bar{x},i})^2 + (v_{\bar{x},i})^2] d\tau \right)^p v_{\bar{x},i} \right\}_x - |v_i|^{q-2} v_i, \\ & \quad i = 1, 2, \dots, M - 1, \end{aligned} \quad (2.4)$$

$$u_0(t) = u_M(t) = v_0(t) = v_M(t) = 0, \quad (2.5)$$

$$u_i(0) = U_{0,i}, \quad v_i(0) = U_{0,i}, \quad i = 0, 1, \dots, M. \quad (2.6)$$

The following statement takes place.

**Theorem 2.** *If  $0 < p \leq 1$ ,  $q \geq 2$  and the initial-boundary value problem (2.1) - (2.3) has the sufficiently smooth solution  $U = U(x, t)$ ,  $V = V(x, t)$ ,*

then the semi-discrete scheme (2.4) - (2.6) converges and the following estimate is true

$$\|u(t) - U(t)\|_h + \|v(t) - V(t)\|_h \leq Ch.$$

Here  $\|\cdot\|_h$  is a discrete analog of the norm of the space  $L_2(0, 1)$  and  $C$  is a positive constant independent of  $h$ .

### 3 Convergence of the Semi-discrete Scheme

In section 2 we constructed Cauchy problem for nonlinear system of ordinary integro-differential equations (2.4) - (2.6) as semi-discrete analog for problem (2.1) - (2.3). The aim of the present section is the proof of the Theorem 2.

Introduce inner products and norms:

$$(r, g)_h = h \sum_{i=1}^{M-1} r_i g_i, \quad (r, g]_h = h \sum_{i=1}^M r_i g_i,$$

$$\|r\|_h = (r, r)_h^{1/2}, \quad \|r\|]_h = (r, r]_h^{1/2}, \quad \|r\|_{q,h}^q = h \sum_{i=1}^{M-1} |r_i|^q.$$

After multiplying scalarly corresponding equations in system (2.4) by  $u(t) = (u_1(t), u_2(t), \dots, u_{M-1}(t))$  and  $v(t) = (v_1(t), v_2(t), \dots, v_{M-1}(t))$  and using discrete analog of integrating by part we get:

$$\frac{d}{dt} \|u(t)\|_h^2 + h \sum_{i=1}^M \left( 1 + \int_0^t [(u_{\bar{x},i})^2 + (v_{\bar{x},i})^2] d\tau \right)^p (u_{\bar{x},i})^2 + \|u(t)\|_{q,h}^q = 0,$$

$$\frac{d}{dt} \|v(t)\|_h^2 + h \sum_{i=1}^M \left( 1 + \int_0^t [(u_{\bar{x},i})^2 + (v_{\bar{x},i})^2] d\tau \right)^p (v_{\bar{x},i})^2 + \|v(t)\|_{q,h}^q = 0.$$

From these relations we obtain the following inequalities:

$$\|u(t)\|_h^2 + \int_0^t \|u_{\bar{x}}\|_h^2 d\tau + \int_0^t \|u(t)\|_{q,h}^q d\tau \leq C,$$

$$\|v(t)\|_h^2 + \int_0^t \|v_{\bar{x}}\|_h^2 d\tau + \int_0^t \|v(t)\|_{q,h}^q d\tau \leq C. \tag{3.7}$$

The a priori estimates (3.7) guarantee the global solvability of the problem (2.4) - (2.6).

**Proof of Theorem 2.** For  $U = U(x, t)$  and  $V = V(x, t)$  we have:

$$\begin{aligned} \frac{dU_i}{dt} - \left\{ \left( 1 + \int_0^t [(U_{\bar{x},i})^2 + (V_{\bar{x},i})^2] d\tau \right)^p U_{\bar{x},i} \right\}_x + |U_i|^{q-2} U_i \\ = \psi_{1,i}(t), \\ \frac{dV_i}{dt} - \left\{ \left( 1 + \int_0^t [(U_{\bar{x},i})^2 + (V_{\bar{x},i})^2] d\tau \right)^p V_{\bar{x},i} \right\}_x + |V_i|^{q-2} V_i \\ = \psi_{2,i}(t), \end{aligned} \quad (3.8)$$

$$i = 1, 2, \dots, M-1,$$

$$U_0(t) = U_M(t) = V_0(t) = V_M(t) = 0, \quad (3.9)$$

$$U_i(0) = U_{0,i}, \quad V_i(0) = V_{0,i}, \quad i = 0, 1, \dots, M, \quad (3.10)$$

where

$$\psi_{k,i}(t) = O(h), \quad k = 1, 2.$$

Let  $z_i(t) = u_i(t) - U_i(t)$  and  $w_i(t) = v_i(t) - V_i(t)$ . From (2.4) - (2.6) and (3.8) - (3.10) we have:

$$\begin{aligned} \frac{dz_i}{dt} - \left\{ \left( 1 + \int_0^t [(u_{\bar{x},i})^2 + (v_{\bar{x},i})^2] d\tau \right)^p u_{\bar{x},i} \right. \\ \left. - \left( 1 + \int_0^t [(U_{\bar{x},i})^2 + (V_{\bar{x},i})^2] d\tau \right)^p U_{\bar{x},i} \right\}_x \\ + |u_i|^{q-2} u_i - |U_i|^{q-2} U_i = -\psi_{1,i}(t), \\ \frac{dw_i}{dt} - \left\{ \left( 1 + \int_0^t [(u_{\bar{x},i})^2 + (v_{\bar{x},i})^2] d\tau \right)^p v_{\bar{x},i} \right. \\ \left. - \left( 1 + \int_0^t [(U_{\bar{x},i})^2 + (V_{\bar{x},i})^2] d\tau \right)^p V_{\bar{x},i} \right\}_x \\ + |v_i|^{q-2} v_i - |V_i|^{q-2} V_i = -\psi_{2,i}(t), \\ z_0(t) = z_M(t) = w_0(t) = w_M(t) = 0, \end{aligned} \quad (3.11)$$

$$z_i(0) = w_i(0) = 0.$$

Multiplying scalarly on  $z(t) = (z_1(t), z_2(t), \dots, z_{M-1}(t))$  the first equation of system (3.11), using the discrete analogue of the formula of integration by parts we get

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} \|z\|^2 + h \sum_{i=1}^M \left\{ \left( 1 + \int_0^t [(u_{\bar{x},i})^2 + (v_{\bar{x},i})^2] d\tau \right)^p u_{\bar{x},i} \right. \\ & \quad \left. - \left( 1 + \int_0^t [(U_{\bar{x},i})^2 + (V_{\bar{x},i})^2] d\tau \right)^p U_{\bar{x},i} \right\} z_{\bar{x},i} \\ & + h \sum_{i=1}^{M-1} (|u_i|^{q-2} u_i - |U_i|^{q-2} U_i) (u_i - U_i) = -h \sum_{i=1}^{M-1} \psi_{1,i} z_i. \end{aligned}$$

Analogously,

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} \|w\|^2 + h \sum_{i=1}^M \left\{ \left( 1 + \int_0^t [(u_{\bar{x},i})^2 + (v_{\bar{x},i})^2] d\tau \right)^p v_{\bar{x},i} \right. \\ & \quad \left. - \left( 1 + \int_0^t [(U_{\bar{x},i})^2 + (V_{\bar{x},i})^2] d\tau \right)^p V_{\bar{x},i} \right\} w_{\bar{x},i} \\ & + h \sum_{i=1}^{M-1} (|v_i|^{q-2} v_i - |V_i|^{q-2} V_i) (v_i - V_i) = -h \sum_{i=1}^{M-1} \psi_{2,i} w_i. \end{aligned}$$

Using monotonicity of the function  $f(r) = |r|^{q-2}r$ , from these two equalities we have

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} (\|z\|^2 + \|w\|^2) + h \sum_{i=1}^M \left\{ \left( 1 + \int_0^t [(u_{\bar{x},i})^2 + (v_{\bar{x},i})^2] d\tau \right)^p u_{\bar{x},i} \right. \\ & \quad \left. - \left( 1 + \int_0^t [(U_{\bar{x},i})^2 + (V_{\bar{x},i})^2] d\tau \right)^p U_{\bar{x},i} \right\} z_{\bar{x},i} \\ & + h \sum_{i=1}^M \left\{ \left( 1 + \int_0^t [(u_{\bar{x},i})^2 + (v_{\bar{x},i})^2] d\tau \right)^p v_{\bar{x},i} \right. \end{aligned} \tag{3.12}$$

$$\begin{aligned}
& - \left( 1 + \int_0^t [(U_{\bar{x},i})^2 + (V_{\bar{x},i})^2] d\tau \right)^p V_{\bar{x},i} \Big\} w_{\bar{x},i} \\
& \leq -h \sum_{i=1}^{M-1} (\psi_{1,i} z_i + \psi_{2,i} w_i).
\end{aligned}$$

Note that,

$$\begin{aligned}
& \left\{ \left( 1 + \int_0^t [(u_{\bar{x},i})^2 + (v_{\bar{x},i})^2] d\tau \right)^p u_{\bar{x},i} \right. \\
& - \left. \left( 1 + \int_0^t [(U_{\bar{x},i})^2 + (V_{\bar{x},i})^2] d\tau \right)^p U_{\bar{x},i} \right\} (u_{\bar{x},i} - U_{\bar{x},i}) \\
& + \left\{ \left( 1 + \int_0^t [(u_{\bar{x},i})^2 + (v_{\bar{x},i})^2] d\tau \right)^p v_{\bar{x},i} \right. \\
& - \left. \left( 1 + \int_0^t [(U_{\bar{x},i})^2 + (V_{\bar{x},i})^2] d\tau \right)^p V_{\bar{x},i} \right\} (v_{\bar{x},i} - V_{\bar{x},i}) \\
& = \int_0^1 \frac{d}{d\xi} \left( 1 + \int_0^t \left\{ [U_{\bar{x},i} + \xi(u_{\bar{x},i} - U_{\bar{x},i})]^2 + [V_{\bar{x},i} + \xi(v_{\bar{x},i} - V_{\bar{x},i})]^2 \right\} d\tau \right)^p \\
& \quad \times [U_{\bar{x},i} + \xi(u_{\bar{x},i} - U_{\bar{x},i})] d\xi (u_{\bar{x},i} - U_{\bar{x},i}) \\
& + \int_0^1 \frac{d}{d\xi} \left( 1 + \int_0^t \left\{ [U_{\bar{x},i} + \xi(u_{\bar{x},i} - U_{\bar{x},i})]^2 + [V_{\bar{x},i} + \xi(v_{\bar{x},i} - V_{\bar{x},i})]^2 \right\} d\tau \right)^p \\
& \quad \times [V_{\bar{x},i} + \xi(v_{\bar{x},i} - V_{\bar{x},i})] d\xi (v_{\bar{x},i} - V_{\bar{x},i}) \\
& = 2p \int_0^1 \left( 1 + \int_0^t \left\{ [U_{\bar{x},i} + \xi(u_{\bar{x},i} - U_{\bar{x},i})]^2 + [V_{\bar{x},i} + \xi(v_{\bar{x},i} - V_{\bar{x},i})]^2 \right\} d\tau \right)^{p-1} \\
& \times \int_0^t \left\{ [U_{\bar{x},i} + \xi(u_{\bar{x},i} - U_{\bar{x},i})] (u_{\bar{x},i} - U_{\bar{x},i}) + [V_{\bar{x},i} + \xi(v_{\bar{x},i} - V_{\bar{x},i})] (v_{\bar{x},i} - V_{\bar{x},i}) \right\} d\tau \\
& \quad \times [U_{\bar{x},i} + \xi(u_{\bar{x},i} - U_{\bar{x},i})] d\xi (u_{\bar{x},i} - U_{\bar{x},i}) \\
& + \int_0^1 \left( 1 + \int_0^t \left\{ [U_{\bar{x},i} + \xi(u_{\bar{x},i} - U_{\bar{x},i})]^2 + [V_{\bar{x},i} + \xi(v_{\bar{x},i} - V_{\bar{x},i})]^2 \right\} d\tau \right)^p
\end{aligned}$$

$$\begin{aligned}
 & \times (u_{\bar{x},i} - U_{\bar{x},i}) d\xi (u_{\bar{x},i} - U_{\bar{x},i}) \\
 & + 2p \int_0^1 \left( 1 + \int_0^t \left\{ [U_{\bar{x},i} + \xi(u_{\bar{x},i} - U_{\bar{x},i})]^2 + [V_{\bar{x},i} + \xi(v_{\bar{x},i} - V_{\bar{x},i})]^2 \right\} d\tau \right)^{p-1} \\
 & \times \int_0^t \left\{ [U_{\bar{x},i} + \xi(u_{\bar{x},i} - U_{\bar{x},i})] (u_{\bar{x},i} - U_{\bar{x},i}) + [V_{\bar{x},i} + \xi(v_{\bar{x},i} - V_{\bar{x},i})] (v_{\bar{x},i} - V_{\bar{x},i}) \right\} d\tau \\
 & \quad \times [V_{\bar{x},i} + \xi(v_{\bar{x},i} - V_{\bar{x},i})] d\xi (v_{\bar{x},i} - V_{\bar{x},i}) \\
 & + \int_0^1 \left( 1 + \int_0^t \left\{ [U_{\bar{x},i} + \xi(u_{\bar{x},i} - U_{\bar{x},i})]^2 + [V_{\bar{x},i} + \xi(v_{\bar{x},i} - V_{\bar{x},i})]^2 \right\} d\tau \right)^p \\
 & \quad \times (v_{\bar{x},i} - V_{\bar{x},i}) d\xi (v_{\bar{x},i} - V_{\bar{x},i}) \\
 & = 2p \int_0^1 \left( 1 + \int_0^t \left\{ [U_{\bar{x},i} + \xi(u_{\bar{x},i} - U_{\bar{x},i})]^2 + [V_{\bar{x},i} + \xi(v_{\bar{x},i} - V_{\bar{x},i})]^2 \right\} d\tau \right)^{p-1} \\
 & \times \int_0^t \left\{ [U_{\bar{x},i} + \xi(u_{\bar{x},i} - U_{\bar{x},i})] (u_{\bar{x},i} - U_{\bar{x},i}) + [V_{\bar{x},i} + \xi(v_{\bar{x},i} - V_{\bar{x},i})] (v_{\bar{x},i} - V_{\bar{x},i}) \right\} d\tau \\
 & \times \left\{ [U_{\bar{x},i} + \xi(u_{\bar{x},i} - U_{\bar{x},i})] (u_{\bar{x},i} - U_{\bar{x},i}) + [V_{\bar{x},i} + \xi(v_{\bar{x},i} - V_{\bar{x},i})] d\xi (v_{\bar{x},i} - V_{\bar{x},i}) \right\} d\xi \\
 & + \int_0^1 \left( 1 + \int_0^t \left\{ [U_{\bar{x},i} + \xi(u_{\bar{x},i} - U_{\bar{x},i})]^2 + [V_{\bar{x},i} + \xi(v_{\bar{x},i} - V_{\bar{x},i})]^2 \right\} d\tau \right)^p \\
 & \quad \times \left[ (u_{\bar{x},i} - U_{\bar{x},i})^2 + (v_{\bar{x},i} - V_{\bar{x},i})^2 \right] d\xi \\
 & = p \int_0^1 \left( 1 + \int_0^t \left\{ [U_{\bar{x},i} + \xi(u_{\bar{x},i} - U_{\bar{x},i})]^2 + [V_{\bar{x},i} + \xi(v_{\bar{x},i} - V_{\bar{x},i})]^2 \right\} d\tau \right)^{p-1} \\
 & \quad \times \frac{d}{dt} \left( \int_0^t \left\{ [U_{\bar{x},i} + \xi(u_{\bar{x},i} - U_{\bar{x},i})] (u_{\bar{x},i} - U_{\bar{x},i}) \right. \right. \\
 & \quad \left. \left. + [V_{\bar{x},i} + \xi(v_{\bar{x},i} - V_{\bar{x},i})] (v_{\bar{x},i} - V_{\bar{x},i}) \right\} d\tau \right)^2 d\xi \\
 & + \int_0^1 \left( 1 + \int_0^t \left\{ [U_{\bar{x},i} + \xi(u_{\bar{x},i} - U_{\bar{x},i})]^2 + [V_{\bar{x},i} + \xi(v_{\bar{x},i} - V_{\bar{x},i})]^2 \right\} d\tau \right)^p
 \end{aligned}$$



$$\times \left[ (u_{\bar{x},i} - U_{\bar{x},i})^2 + (v_{\bar{x},i} - V_{\bar{x},i})^2 \right] d\xi.$$

After substituting this equality in (3.12), integrating received equality on  $(0, t)$  and using formula of integrating by parts we get

$$\begin{aligned} & \|z\|^2 + \|w\|^2 + 2h \sum_{i=1}^M \int_0^t \int_0^1 \left( 1 + \int_0^{t'} \left\{ [U_{\bar{x},i} + \xi(u_{\bar{x},i} - U_{\bar{x},i})]^2 \right. \right. \\ & \left. \left. + [V_{\bar{x},i} + \xi(v_{\bar{x},i} - V_{\bar{x},i})]^2 \right\} d\tau' \right)^p \left[ (u_{\bar{x},i} - U_{\bar{x},i})^2 + (v_{\bar{x},i} - V_{\bar{x},i})^2 \right] d\xi d\tau \\ & + 2ph \sum_{i=1}^M \int_0^1 \left( 1 + \int_0^t \left\{ [U_{\bar{x},i} + \xi(u_{\bar{x},i} - U_{\bar{x},i})]^2 + [V_{\bar{x},i} + \xi(v_{\bar{x},i} - V_{\bar{x},i})]^2 \right\} d\tau \right)^{p-1} \\ & \quad \times \left( \int_0^t \left\{ [U_{\bar{x},i} + \xi(u_{\bar{x},i} - U_{\bar{x},i})] (u_{\bar{x},i} - U_{\bar{x},i}) \right. \right. \\ & \quad \left. \left. + [V_{\bar{x},i} + \xi(v_{\bar{x},i} - V_{\bar{x},i})] (v_{\bar{x},i} - V_{\bar{x},i}) \right\} d\tau \right)^2 d\xi \\ & - 2p(p-1)h \sum_{i=1}^M \int_0^1 \int_0^t \left( 1 + \int_0^{t'} \left\{ [U_{\bar{x},i} + \xi(u_{\bar{x},i} - U_{\bar{x},i})]^2 \right. \right. \\ & \quad \left. \left. + [V_{\bar{x},i} + \xi(v_{\bar{x},i} - V_{\bar{x},i})]^2 \right\} d\tau' \right)^{p-2} \\ & \quad \times \left\{ [U_{\bar{x},i} + \xi(u_{\bar{x},i} - U_{\bar{x},i})]^2 + [V_{\bar{x},i} + \xi(v_{\bar{x},i} - V_{\bar{x},i})]^2 \right\} \\ & \quad \times \left( \int_0^{t'} \left\{ [U_{\bar{x},i} + \xi(u_{\bar{x},i} - U_{\bar{x},i})] (u_{\bar{x},i} - U_{\bar{x},i}) \right. \right. \\ & \quad \left. \left. + [V_{\bar{x},i} + \xi(v_{\bar{x},i} - V_{\bar{x},i})] (v_{\bar{x},i} - V_{\bar{x},i}) \right\} d\tau' \right)^2 d\xi d\tau \\ & = -2h \sum_{i=1}^{M-1} \int_0^t (\psi_{1,i} z_i + \psi_{2,i} w_i) d\tau. \end{aligned}$$

Taking into account relation  $0 < p \leq 1$  from last equality we have

$$\begin{aligned} \|z(t)\|_h^2 + \|w(t)\|_h^2 &\leq \int_0^t (\|z(\tau)\|_h^2 + \|w(\tau)\|_h^2) d\tau \\ &+ \int_0^t (\|\psi_1\|_h^2 + \|\psi_2\|_h^2) d\tau. \end{aligned} \quad (3.13)$$

From (3.13) using Gronwall's inequality we get validity of the Theorem 2.

Note that investigated semi-discrete scheme (2.4) - (2.6) is using for numerical solution of the problem (2.1) - (2.3) by natural discretisation of time derivative and integral as it are given for example in [23], [24] for the case  $p = 1$ . Solving the obtaining finite difference scheme we use a algorithm analogical to [25]. So, it is necessary to use Newton iterative process [27]. According to this method the great numbers of numerical experiments are carried out. These experiments agree with the theoretical results given in the Theorems 1 and 2.

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