# ON ONE VERSION OF THE CHARACTERISTIC PROBLEM FOR A NONLINEAR OSCILLATION EQUATION 

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## Abstract

The paper proposes one nonlinear version of the characteristic problem for a nonlinear oscillation equation, which makes possible to simultaneously define regular solutions and its extension domains. The conditions of the problem are set forth on non interesting arcs of curves characteristic to various families.

Key words and phrases: Characteristics, characteristic invariants, general integral, solution definition domain.

AMS subject classification: 35L70.

## 1 Introduction

In the work [1] dedicated to linear non-strictly hyperbolic equations much attention is given to the well-known second order equation of nonlinear oscillations

$$
u_{y}^{4} u_{x x}-u_{y y}=0
$$

which depending on the behavior and values of the first order derivative $u_{y}$ of the sought solution $u(x, y)$ may generate parabolically. Its general integral is represented by means of the generator of groups of solutions which are based on contact transformations. However the question as to the solvability of problems was not considered. If the method proposed there is used for construction of a general integral for an equation with the none zero right-hand part, then we are faced with great and even insurmountable difficulties.

In the present paper, an analogous question is considered for the equation

$$
\begin{equation*}
u_{y}^{4} u_{x x}-u_{y y}=c x^{-2} u u_{y}^{4}, \quad c=\text { const }, \tag{1}
\end{equation*}
$$

which is interesting not only from the theoretical viewpoint, but also as having various practical applications [2-5].

Like in the case for general nonlinear hyperbolic equations, for equation (1), too, linear formulations of boundary value or characteristic problems,
except for the Cauchy initial boundary value problem, are meaningless. This fact is caused by the dependence of characteristic families on yet unknown solutions.

## 2 General Integral

The right-hand part of equation (1) along the ordinate axis is unbounded. This property makes it possible to attribute this equation to the class EulerDarboux e quations [6-8]. By performing multiplication by a factor defining this unboundedness and making some assumptions, we can consider, instead of (1), an equation with the degeneration not only of hypebolicity but of order, too. However, there may exist solutions along which equation (1) remains hyperbolic. In other words, these are solutions for which the characteristic directions defined by the characteristic roots

$$
\lambda_{1}=u_{y}^{-2}, \quad \lambda_{2}=-u_{y}^{-2}
$$

do not coincide anywhere. Naturally, the class of hyperbolic solutions of equation (1) is defined by the condition

$$
0 \neq\left|u_{y}(x, y)\right|<\infty
$$

Before formulating a nonlinear analogue of the Goursat problem, let us clarify the general properties of hyperbolic solutions of equation (1). For this we will use the classical characteristic method [9].

As is known, the characteristic roots $\lambda_{1}, \lambda_{2}$ give the differential relations of characteristic directions

$$
\begin{equation*}
u_{y}^{2} d y-d x=0, \quad u_{y}^{2} d y+d x=0 . \tag{2}
\end{equation*}
$$

If equation (1) is considered with (2) taken into account, then we come to the following differential characteristic relations

$$
x^{2} u_{y}^{4} d u_{x}-x^{2} u_{y}^{2} d u_{y}-c u_{y}^{4} u d x=0, \quad x^{2} u_{y}^{4} d u_{x}+x^{2} u_{y}^{2} d u_{y}-c u_{y}^{4} u d x=0 .
$$

Assuming that $c>-\frac{1}{4}$, each of the characteristic systems of equation (1) admits exactly two first integrals and they are defined in the explicit form [10]

$$
\left\{\begin{array}{l}
\xi \equiv\left(u_{y}^{-1}+u_{x}\right) x^{\alpha}-\alpha u x^{\alpha-1}  \tag{3}\\
\xi_{1} \equiv\left(u_{y}^{-1}+u_{x}\right) x^{1-\alpha}-(1-\alpha) u x^{-\alpha}
\end{array}\right.
$$

for the family of the root $\lambda_{1}$ and

$$
\left\{\begin{array}{l}
\eta \equiv\left(u_{y}^{-1}-u_{x}\right) x^{\alpha}+\alpha u x^{\alpha-1}  \tag{4}\\
\eta_{1} \equiv\left(u_{y}^{-1}-u_{x}\right) x^{1-\alpha}+(1-\alpha) u x^{-\alpha}
\end{array} \quad, \quad \alpha=\frac{1}{2}(1+\sqrt{4 c+1})\right.
$$

for the family of the root $\lambda_{2}$.
Due to these two pairs of first integrals $\left(\xi, \xi_{1}\right)$ and $\left(\eta, \eta_{1}\right)$ which are actually characteristic invariants, it follows that in the class of hyperbolic solutions we can construct two intermediate integrals

$$
\xi_{1}=\varphi^{\prime}(\xi), \quad \eta_{1}=\psi^{\prime}(\eta)
$$

of equation (1). In these integrals we denote by $\varphi, \psi$ arbitrary functions smooth enough to ensure the differentiability of the sought solution up to second order.

If $\phi, \psi \in c^{3}\left(R_{1}\right)$, then equation (1) is equivalent to a triple of the following relations [10]

$$
\begin{align*}
x & =\left(\frac{\varphi^{\prime}(\xi)+\psi^{\prime}(\eta)}{\xi+\eta}\right)^{\frac{1}{1-2 \alpha}} ;  \tag{5}\\
y & =\frac{1}{4(1-2 \alpha)}\left[(\xi+\eta)\left(\psi^{\prime}(\eta)-\varphi^{\prime}(\xi)\right)+2(\phi(\xi)-\psi(\eta)] ;\right.  \tag{6}\\
u & =\frac{1}{1-2 \alpha}\left[\xi\left(\frac{\varphi^{\prime}(\xi)+\psi^{\prime}(\eta)}{\xi+\eta}\right)^{\frac{1-\alpha}{1-2 \alpha}}-\varphi^{\prime}(\xi)\left(\frac{\varphi^{\prime}(\xi)+\psi^{\prime}(\eta)}{\xi+\eta}\right)^{\frac{\alpha}{1-2 \alpha}}\right], \tag{7}
\end{align*}
$$

which in a certain sense can be taken as a general integral of equation (1), whereas the invariants $\xi, \eta$ play the role of characteristic variables.

## 3 Statement of the Problem and the Method of it Solve

Suppose we are given two arcs $\gamma_{1}, \gamma_{2}$ drawn from the common points $\left(a, f_{1}(a)\right),\left(a, f_{2}(a)\right), f_{2}(a)<f_{1}(a)$ and let them be given in the explicit form

$$
\gamma_{1}: y=f_{1}(x), \quad a \leq x \leq b, \quad a>0
$$

and

$$
\begin{equation*}
\gamma_{2}: y=f_{2}(x), a \leq x \leq c . \tag{8}
\end{equation*}
$$

Assume that the functions $f_{1}$ and $f_{2}$ are three times continuously differentiable and the arc $\gamma_{1}$ monotonically ascends, whereas the arc $\gamma_{2}$, vice versa, monotonically descends.

The characteristic problem. Find a regular hyperbolic solution $u(x, y)$ of equation (1) and, simultaneously with it, a domain of its ex-
tension when the curves $\gamma_{1}$ and $\gamma_{2}$ are the arcs of the characteristics, and

$$
\begin{align*}
u\left(a, f_{1}(a)\right) & =\vartheta_{1},  \tag{9}\\
u_{x}\left(a, f_{1}(a)\right) & =\delta_{1},  \tag{10}\\
u\left(a, f_{2}(a)\right) & =\vartheta_{2},  \tag{11}\\
u_{x}\left(a, f_{2}(a)\right) & =\delta_{2} . \tag{12}
\end{align*}
$$

By the conditions of the problem we have

$$
f_{1}^{\prime}(x)=u_{y}^{-2}, \quad f_{2}^{\prime}(x)=-u_{y}^{-2}
$$

which define the values of the derivative $u_{y}$ on curves $\gamma_{1}$ and $\gamma_{2}$. But they are defined non-uniquely:

$$
\begin{equation*}
\left.u_{y}\right|_{\gamma_{1}}= \pm \frac{1}{\sqrt{f_{1}^{\prime}(x)}},\left.\quad u_{y}\right|_{\gamma_{2}}= \pm \frac{1}{\sqrt{-f_{2}^{\prime}(x)}} \tag{13}
\end{equation*}
$$

To obtain the continuity of the derivative, we should take the right-hand parts of the latter equalities of the same sign.

Let us first consider the case

$$
u_{y}\left(a, f_{1}(a)\right)=\frac{1}{\sqrt{f_{1}^{\prime}(a)}}
$$

Using (9)-(13) we calculate in a straightforward manner the values of characteristic invariants $\xi, \xi_{1}$ at the point $\left(a, f_{1}(a)\right)$ and characteristic invariants $\eta, \eta_{1}$ at the point $\left(a, f_{2}(a)\right)$. Introduce for them the notation

$$
\left.\xi\right|_{\left(a, f_{1}(a)\right)} \equiv \xi^{[a]},\left.\quad \xi_{1}\right|_{\left(a, f_{1}(a)\right)} \equiv \xi_{1}^{[a]},\left.\quad \eta\right|_{\left(a, f_{2}(a)\right)} \equiv \eta^{[a]},\left.\quad \eta_{1}\right|_{\left(a, f_{2}(a)\right)} \equiv \eta_{1}^{[a]}
$$

By the property of characteristic invariants that they are constant along the respective curves, we have

$$
\begin{equation*}
\left.\xi\right|_{\gamma_{1}} \equiv \xi^{[a]},\left.\quad \xi_{1}\right|_{\gamma_{1}} \equiv \xi_{1}^{[a]} \tag{14}
\end{equation*}
$$

and

$$
\begin{equation*}
\left.\eta\right|_{\gamma_{2}} \equiv \eta^{[a]},\left.\quad \eta_{1}\right|_{\gamma_{2}} \equiv \eta_{1}^{[a]} \tag{15}
\end{equation*}
$$

Let us now turn to the general integral (5)-(7) defined in terms of the characteristic invariants $\xi, \eta$ and, using (5), (7), establish the relation between the values $x, u$ along the arc $\gamma_{1}$ :

$$
\begin{equation*}
\left.u\right|_{\gamma_{1}}=\frac{1}{1-2 \alpha}\left[\xi^{[a]} x^{1-\alpha}-\xi_{1}^{[a]} x^{\alpha}\right] . \tag{16}
\end{equation*}
$$

From the second relation (14) we obtain

$$
\begin{equation*}
\left.u_{x}\right|_{\gamma_{1}}=\frac{1-\alpha}{1-2 \alpha} \xi^{[a]} x^{-\alpha}-\frac{\alpha}{1-2 \alpha} \xi_{1}^{[a]} x^{\alpha-1}-\sqrt{f_{1}^{\prime}(x)} . \tag{17}
\end{equation*}
$$

By (4), (15) we have

$$
\begin{equation*}
\left.u\right|_{\gamma_{2}}=\frac{1}{2 \alpha-1}\left[\eta^{[a]} x^{1-\alpha}-\eta_{1}^{[a]} x^{\alpha}\right] . \tag{18}
\end{equation*}
$$

A pair of relations (15), (18) enables us to define the relation of the derivative $u_{x}$ with the variable $x$ along the arc $\gamma_{2}$ as follows

$$
\begin{equation*}
\left.u_{x}\right|_{\gamma_{2}}=\sqrt{-f_{2}^{\prime}(x)}+\frac{1-\alpha}{2 \alpha-1} \eta^{[a]} x^{-\alpha}-\frac{\alpha}{2 \alpha-1} \eta_{1}^{[a]} x^{\alpha-1} . \tag{19}
\end{equation*}
$$

Thus the conditions of problem (1), (9)-(12) make it possible to define along both characteristic arcs $\gamma_{1}, \gamma_{2}$ the interdependence between the argument $x$, the solution $u(x, y)$ of equation (1) and its first order derivatives. But the aim we pursue is to define the solution not only on the characteristics but also outside the data carrier. To this end, we will again try to use the general properties of characteristic invariants. On the arc $\gamma_{1}$, we can represent the invariants $\eta, \eta_{1}$ from the other family as functions of the argument $x$. This is done by substitution of (13), (16), (17) into (4). On the arc $\gamma_{2}$, the values $\xi_{1}, \xi_{2}$ are constructed in the same manner.

From an arbitrary point $P_{2}\left(x_{2}, f\left(x_{2}\right)\right) \in \gamma_{2}$ let us draw the characteristic $\gamma_{3}$ of the family of the root $\lambda_{1}$. The constants along will be the invariants $\xi$ and $\xi_{1}$. Another point $P_{1}\left(x_{1}, f_{1}\left(x_{1}\right)\right) \in \gamma_{1}$ is treated analogously and the characteristic of the family of the root $\lambda_{2}$ drawn from it is denoted by $\gamma_{4}$. The values of the invariants $\eta, \eta_{1}$ at the point $P_{1}$ are preserved all along the arc $\gamma_{4}$. According to the general theory, a set of points of intersection of the analogous characteristics $\gamma_{3}, \gamma_{4}$ defines the domain of solution extension. If there exits a point of intersection of the $\operatorname{arcs} \gamma_{3}, \gamma_{4}$, we denote it by $M\left(x^{0}, y^{0}\right)$.

Let us introduce the notation

$$
\left.\xi\right|_{P_{2}} \equiv \xi^{\left[x_{2}\right]},\left.\quad \xi_{1}\right|_{P_{2}} \equiv \xi_{1}^{\left[x_{2}\right]},\left.\quad \eta\right|_{P_{1}} \equiv \eta^{\left[x_{1}\right]},\left.\quad \eta_{1}\right|_{P_{1}} \equiv \eta_{1}^{\left[x_{1}\right]},
$$

As it has already been said, the following relations hold true

$$
\begin{array}{ll}
\left.\xi\right|_{\gamma_{3}}=\xi^{\left[x_{2}\right]}, & \left.\xi_{1}\right|_{\gamma_{3}}=\xi_{1}^{\left[x_{2}\right]}, \\
\left.\eta\right|_{\gamma_{4}}=\eta^{\left[x_{1}\right]}, & \left.\eta_{1}\right|_{\gamma_{4}}=\eta_{1}^{\left[x_{1}\right]} \tag{21}
\end{array}
$$

and they are simultaneously fulfilled at the point $M\left(x^{0}, y^{0}\right)$ of intersection of the arcs $\gamma_{3}, \gamma_{4}$.

Thus we obtain the following system with respect to $x^{0}, u^{0}=u\left(x^{0}, y^{0}\right)$, $u_{x}^{0}=u_{x}\left(x^{0}, y^{0}\right), u_{y}^{0}=u_{y}\left(x^{0}, y^{0}\right):$

$$
\left\{\begin{array}{l}
\left(u_{y}^{0-1}+u_{x}^{0}\right) x^{0^{\alpha}}-\alpha u^{0} x^{0^{\alpha-1}}=\xi^{\left[x_{2}\right]}  \tag{22}\\
\left(u_{y}^{0-1}+u_{x}^{0}\right) x^{0^{1-\alpha}}-(1-\alpha) u^{0} x^{0^{-\alpha}}=\xi_{1}^{\left[x_{2}\right]} \\
\left(u_{y}^{0-1}-u_{x}^{0}\right) x^{0^{\alpha}}+\alpha u^{0} x^{0^{\alpha-1}}=\eta^{\left[x_{1}\right]} \\
\left(u_{y}^{0-1}-u_{x}^{0}\right) x^{0^{1-\alpha}}-(1-\alpha) u^{0} x^{0^{-\alpha}}=\eta_{1}^{\left[x_{1}\right]}
\end{array}\right.
$$

We have already dealt with a system of form (22) when constructing the general integral (5)-(7). Hence to solve system (22), we perform analogous actions with the only difference that (22) will be considered not with respect to the functions $x, u, u_{x}, u_{y}$ but with respect to their concrete values at the point $M\left(x^{0}, y^{0}\right)$. The right-hand parts are also concrete constant values well defined by the conditions of problem (1), (9)-(12).

To find $y^{0}=y\left(x_{1}, x_{2}\right)$ we resort to relation (6). Taking into consideration the explicit equations of the arcs $\gamma_{1}, \gamma_{2}$, we finally define the value $y=y^{0}$ in terms of $x_{1}, x_{2}$. The values $x^{0}, y^{0}, u^{0}$ are expressed as follows

$$
\begin{align*}
x^{0} & =\left(\frac{\xi_{1}^{\left[x_{2}\right]}+\eta_{1}^{\left[x_{1}\right]}}{\xi^{\left[x_{2}\right]}+\eta^{\left[x_{1}\right]}}\right)^{\frac{1}{1-2 \alpha}}  \tag{23}\\
y^{0} & =f_{1}\left(x_{1}\right)+f_{2}\left(x_{2}\right) \\
& +\frac{1}{4(1-2 \alpha)}\left[\xi^{\left[x_{2}\right]} \eta_{1}^{\left[x_{1}\right]}-\eta^{\left[x_{1}\right]} \xi_{1}^{\left[x_{2}\right]}-\xi^{[a]} \eta_{1}^{\left[x_{1}\right]}+\eta^{\left[x_{1}\right]} \xi_{1}^{[a]}\right. \\
& \left.-\xi^{\left[x_{2}\right]} \eta_{1}^{[a]}+\xi_{1}^{\left[x_{2}\right]} \eta^{[a]}+\xi^{[a]} \eta_{1}^{[a]}-\eta^{[a]} \xi_{1}^{[a]}\right]  \tag{24}\\
u^{0} & =\frac{1}{1-2 \alpha}\left[\xi^{\left[x_{2}\right]}\left(\frac{\xi_{1}^{\left[x_{2}\right]}+\eta_{1}^{\left[x_{1}\right]}}{\xi^{\left[x_{2}\right]}+\eta^{\left[x_{1}\right]}}\right)^{\frac{1-\alpha}{1-2 \alpha}}-\xi_{1}^{\left[x_{2}\right]}\left(\frac{\xi_{1}^{\left[x_{2}\right]}+\eta_{1}^{\left[x_{1}\right]}}{\xi^{\left[x_{2}\right]}+\eta^{\left[x_{1}\right]}}\right)^{\frac{\alpha}{1-2 \alpha}}\right] . \tag{25}
\end{align*}
$$

We have obtained the definition domain $D$ of the solution $u\left(x^{0}, y^{0}\right)$ of problem (1), (9)-(12) for the current $x_{1}, x_{2}$. This domain is well defined by relations (23), (24), where the expressions of $x^{0}, y^{0}$ are presented depending on $x_{1}, x_{2}$. We consider the values of these functions as the current coordinates describing the domain $D$. Then, taking into consideration the expressions of invariants along the arcs $\gamma_{3}, \gamma_{4}$ and their values at $(a, 0)$, we can represent the solution of the considered problem (1), (9)-(12) by the following three equalities

$$
\begin{gathered}
x=F_{1}\left(x_{1}, x_{2}\right) \\
\equiv\left(2 \sqrt{-f_{2}^{\prime}\left(x_{2}\right)} x_{2}^{1-\alpha}+2 \sqrt{-f_{1}^{\prime}\left(x_{1}\right)} x_{1}^{1-\alpha}\right.
\end{gathered}
$$

$$
\begin{gather*}
\begin{array}{c}
\left.-\left(\sqrt{f_{1}^{\prime}(a)}+\sqrt{-f_{2}^{\prime}(a)}+\delta_{1}-\delta_{2}\right) a^{1-\alpha}+\left(\vartheta_{1}-\vartheta_{2}\right)(1-\alpha) a^{-\alpha}\right)^{\frac{1}{1-2 \alpha}} \\
\times\left(2 \sqrt{-f_{2}^{\prime}\left(x_{2}\right)} x_{2}^{\alpha}+2 \sqrt{f_{1}^{\prime}\left(x_{1}\right)} x_{1}^{\alpha}\right. \\
\left.-\left(\sqrt{f_{1}^{\prime}(a)}+\sqrt{-f_{2}^{\prime}(a)}+\delta_{1}-\delta_{2}\right) a^{\alpha}\left(\vartheta_{1}-\vartheta_{2}\right) \alpha a^{\alpha-1}\right)^{\frac{1}{2 \alpha-1}} ; \\
y=G_{1}\left(x_{1}, x_{2}\right) \equiv f_{1}\left(x_{1}\right)+f_{2}\left(x_{2}\right) \\
+\frac{1}{4(1-2 \alpha)}\left[4 \sqrt{f_{1}^{\prime}\left(x_{1}\right)} \sqrt{-f_{2}^{\prime}\left(x_{2}\right)\left(x_{1}^{1-\alpha} x_{2}^{\alpha}-x_{1}^{\alpha} x_{2}^{1-\alpha}\right)}\right. \\
+2\left(\sqrt{-f_{2}^{\prime}\left(x_{2}\right)} x_{2}^{1-\alpha}-\sqrt{f_{1}^{\prime}\left(x_{1}\right)} x_{1}^{1-\alpha}\right) \\
\times\left(\left(\sqrt{f_{1}^{\prime}(a)}+\sqrt{-f_{2}^{\prime}(a)}+\delta_{1}-\delta_{2}\right) a^{\alpha}+\alpha\left(\vartheta_{2}-\vartheta_{1}\right) a^{\alpha-1}\right) \\
\quad+2\left(\sqrt{f_{1}^{\prime}\left(x_{1}\right)} x_{1}^{\alpha}-\sqrt{-f_{2}^{\prime}\left(x_{2}\right)} x_{2}^{\alpha}\right) \\
\left.\times\left(\left(\sqrt{f_{1}^{\prime}(a)}+\sqrt{-f_{2}^{\prime}(a)}+\delta_{1}-\delta_{2}\right) a^{1-\alpha}+(1-\alpha)\left(\vartheta_{2}-\vartheta_{1}\right) a^{-\alpha}\right)\right] \\
\times\left[\left(2 \sqrt{-f_{2}^{\prime}\left(x_{2}\right)} x_{2}^{\alpha}-\left(\sqrt{-f_{2}^{\prime}(a)}-\delta_{2}\right) a^{\alpha}-\alpha \vartheta_{2} a^{\alpha-1}\right) F_{1}^{1-\alpha}\left(x_{1}, x_{2}\right)\right. \\
\quad+\left(2 \sqrt{-f_{2}^{\prime}\left(x_{2}\right)} x_{2}^{1-\alpha}\right. \\
\left.\left.\quad+\left(\sqrt{-f_{2}^{\prime}(a)}-\delta_{2}\right) a^{1-\alpha}-(1-\alpha) \vartheta_{2} a^{-\alpha}\right) F_{1}^{\alpha}\left(x_{1}, x_{2}\right)\right]
\end{array}
\end{gather*}
$$

They simultaneously define both the domain $D_{1}$ and the sought solution.

For the characteristics of one and the same family not to intersect each other, it is sufficient that the conditions

$$
\begin{align*}
& {\left[F_{1}\left(x_{1}^{0}, x_{2}\right)-F_{1}\left(x_{1}^{\prime}, x_{2}\right)\right]^{2}+\left[G_{1}\left(x_{1}^{0}, x_{2}\right)-G_{1}\left(x_{1}^{\prime}, x_{2}\right)\right]^{2} \neq 0,}  \tag{29}\\
& {\left[F_{1}\left(x_{1}, x_{2}^{0}\right)-F_{1}\left(x_{1}, x_{2}^{\prime}\right)\right]^{2}+\left[G_{1}\left(x_{1}, x_{2}^{0}\right)-G_{1}\left(x_{1}, x_{2}^{\prime}\right)\right]^{2} \neq 0} \tag{30}
\end{align*}
$$

be fulfilled for any fixed values $x_{1}^{0} \neq x_{1}^{\prime}, x_{1}^{0}, x_{1}^{\prime} \in[a, b]$ and $x_{2}^{0} \neq x_{2}^{\prime}$, $x_{2}^{0}, x_{2}^{\prime} \in[a, c]$ and for the parameters $x_{1} \in[a, b], x_{2} \in[a, c]$.

If

$$
\left.u_{y}\right|_{\gamma_{1}}=-\frac{1}{\sqrt{f_{1}^{\prime}(x)}},\left.\quad u_{y}\right|_{\gamma_{2}}=-\frac{1}{\sqrt{f_{2}^{\prime}(x)}}
$$

then by a reasoning analogous to that in the preceding case we see that problem (1), (9)-(12) has, in addition to (26)-(28), one more solution rep-
resented by the formulas

$$
\begin{gather*}
x=F_{2}\left(x_{1}, x_{2}\right) \equiv\left(2 \sqrt{-f_{2}^{\prime}\left(x_{2}\right)} x_{2}^{1-\alpha}+2 \sqrt{f_{1}^{\prime}\left(x_{1}\right)} x_{1}^{1-\alpha}\right. \\
\left.-\left(\sqrt{f_{1}^{\prime}(a)}+\sqrt{-f_{2}^{\prime}(a)}+\delta_{2}-\delta_{1}\right) a^{1-\alpha}+\left(\vartheta_{2}-\vartheta_{1}\right)(1-\alpha) a^{-\alpha}\right)^{\frac{1}{1-2 \alpha}} \\
\times\left(2 \sqrt{-f_{2}^{\prime}\left(x_{2}\right)} x_{2}^{\alpha}+2 \sqrt{f_{1}^{\prime}\left(x_{1}\right)} x_{1}^{\alpha}\right. \\
\left.-\left(\sqrt{f_{1}^{\prime}(a)}+\sqrt{-f_{2}^{\prime}(a)}+\delta_{2}-\delta_{1}\right) a^{\alpha}+\left(\vartheta_{2}-\vartheta_{1}\right) \alpha a^{\alpha-1}\right)^{\frac{1}{2 \alpha-1}}  \tag{31}\\
y=G_{2}\left(x_{1}, x_{2}\right) \equiv f_{1}\left(x_{1}\right)+f_{2}\left(x_{2}\right) \\
+\frac{1}{4(1-2 \alpha)}\left[4 \sqrt{f_{1}^{\prime}\left(x_{1}\right)} \sqrt{-f_{2}^{\prime}\left(x_{2}\right)}\left(x_{1}^{1-\alpha} x_{2}^{\alpha}-x_{1}^{\alpha} x_{2}^{1-\alpha}\right)\right. \\
-2\left(\sqrt{-f_{2}^{\prime}\left(x_{2}\right)} x_{2}^{1-\alpha}-\sqrt{f_{1}^{\prime}\left(x_{1}\right)} x_{1}^{1-\alpha}\right) \\
\times\left(\left(\delta_{1}-\delta_{2}-\sqrt{f_{1}^{\prime}(a)}-\sqrt{f_{2}^{\prime}(a)}\right) a^{\alpha}+\alpha\left(\vartheta_{2}-\vartheta_{1}\right) a^{\alpha-1}\right) \\
-2\left(\sqrt{f_{1}^{\prime}\left(x_{1}\right)} x_{1}^{\alpha}-\sqrt{-f_{2}^{\prime}\left(x_{2}\right)} x_{2}^{\alpha}\right) \\
\quad u=\frac{1}{2 \alpha-1}  \tag{32}\\
\left.\times\left(\left(\delta_{1}-\delta_{2}-\sqrt{f_{1}^{\prime}(a)}-\sqrt{f_{2}^{\prime}(a)}\right) a^{\alpha}+\alpha\left(\vartheta_{2}-\vartheta_{1}\right) a^{\alpha-1}\right)\right] \\
\times\left[\left(2 \sqrt{-f_{2}^{\prime}\left(x_{2}\right)} x_{2}^{\alpha}-\left(\delta_{2}-\sqrt{-f_{2}^{\prime}(a)}\right) a^{\alpha}+\alpha \vartheta_{2} a^{\alpha-1}\right) F_{2}^{1-\alpha}\left(x_{1}, x_{2}\right)\right. \\
+\left(2 \sqrt{-f_{2}^{\prime}\left(x_{2}\right)} x_{2}^{1-\alpha}-\left(\delta_{2}-\sqrt{-f_{2}^{\prime}(a)}\right) a^{1-\alpha}\right. \\
\left.\left.\quad+(1-\alpha) \vartheta_{2} a^{-\alpha}\right) F_{2}^{\alpha}\left(x_{1}, x_{2}\right)\right] \tag{33}
\end{gather*}
$$

equalities (31), (32) define the definition domain $D_{2}$.
For the characteristics of one and the same family not to intersect each other in the domain $D_{2}$, it is sufficient that the conditions

$$
\begin{align*}
& {\left[F_{2}\left(x_{1}^{0}, x_{2}\right)-F_{2}\left(x_{1}^{\prime}, x_{2}\right)\right]^{2}+\left[G_{2}\left(x_{1}^{0}, x_{2}\right)-G_{2}\left(x_{1}^{\prime}, x_{2}\right)\right]^{2} \neq 0}  \tag{34}\\
& {\left[F_{2}\left(x_{1}, x_{2}^{0}\right)-F_{2}\left(x_{1}, x_{2}^{\prime}\right)\right]^{2}+\left[G_{2}\left(x_{1}, x_{2}^{0}\right)-G_{2}\left(x_{1}, x_{2}^{\prime}\right)\right]^{2} \neq 0} \tag{35}
\end{align*}
$$

be fulfilled for any fixed values $x_{1}^{0} \neq x_{1}^{\prime}, x_{1}^{0}, x_{1}^{\prime} \in[a, b]$ and $x_{2}^{0} \neq x_{2}^{\prime}$, $x_{2}^{0}, x_{2}^{\prime} \in[a, c]$ and for the parameters $x_{1} \in[a, b], x_{2} \in[a, c]$.

The following theorem is true.

Theorem. Under conditions (29), (30), (34), (35) for any real branch of the multi-valued functions (26), (31) there are regular hyperbolic solutions of problem (1), (9)-(12) represented in the explicit form by formulas (26)(28) in the domain $D_{1}$ and (31)-(33) in the domain $D_{2}$. Where $D_{1}$ domain is bounded by arcs of characteristic curves

$$
\begin{aligned}
& \Gamma_{1}: x=F_{1}\left(a, x_{2}\right), \quad y=G_{1}\left(a, x_{2}\right) ; \\
& \Gamma_{2}: x=F_{1}\left(x_{1}, c\right), \quad y=G_{1}\left(x_{1}, c\right) ; \\
& \Gamma_{3}: x=F_{1}\left(b, x_{2}\right), \quad y=G_{1}\left(b, x_{2}\right) ; \\
& \Gamma_{4}: x=F_{1}\left(x_{1}, d\right), \quad y=G_{1}\left(x_{1}, d\right) \text {; }
\end{aligned}
$$

and domain $D_{2}$ is bounded by arcs of characteristic curves

$$
\begin{aligned}
& \Gamma_{5}: x=F_{2}\left(a, x_{2}\right), y=G_{2}\left(a, x_{2}\right) ; \\
& \Gamma_{6}: x=F_{2}\left(x_{1}, c\right), y=G_{2}\left(x_{1}, c\right) ; \\
& \Gamma_{7}: x=F_{2}\left(b, x_{2}\right), y=G_{2}\left(b, x_{2}\right) ; \\
& \Gamma_{8}: x=F_{2}\left(x_{1}, d\right), y=G_{2}\left(x_{1}, d\right) .
\end{aligned}
$$

## References

1. S. Kumei, G. W. Bluman, When nonlinear differential equations are equivalent to linear differential equations, SIAM J. Appl. Math. 42 (1982), No. 5, 1157-1173.
2. J. Hadamard, Leçons sur un propagation de ondes, Paris, 2010.
3. A. V. Bitsadze, Equations of mixed type (Russian), Ergebnisse der Wissenschaft. Physikalisch-mathematische Wissenschaften. 2. Verlag der Akademie der Wissenschaften der UdSSR, Moscow, 1959.
4. L. Bers, Mathematical aspects of subsonic and transonic gas dynam$i c s$, Surveys in Applied Mathematics, Vol. 3 John Wiley \& Sons, Inc., New York; Chapman \& Hall, Ltd., London, 1958.
5. J. Gvazava, Some classes of quasilinear equations of mixed type (Russian), Trudy Tbiliss. Mat. Inst. Razmadze Akad. Nauk Gruzin. SSR 67 (1981), 94 pp.
6. A. V. Bitsadze, Some classes of partial differential equations. Translated from the Russian by H. Zahavi. Advanced Studies in Contemporary Mathematics, 4. Gordon and Breach Science Publishers, New York, 1988
7. E. T. Whittaker, G. N. Watson, A course of modern analysis. I, II (Russian). Translation from the English. 2nd ed. Staatsverlag für physikalisch-mathematische Literatur, 1963.
8. F. Tricomi, Lectures on partial equations. Moscow, 1957.
9. E. Goursat, Leçons sur l'intégration des équations aux dérivées partielles du second ordre à deux variables indépendantes. Tome II: La méthode de Laplace. Les systémes en involutions. La méthode de M. Darboux. Les équations de la première classe. Transformations des équations du second ordre. Généralisations diverses. 2 Notes. A. Hermann, Paris, 1898.
10. R. G.Bitsadze, General representation of the solutions of a quasilinear equation of a problem of nonlinear oscillations (Russian), Soobshch. Akad. Nauk Gruzin. SSR 128 (1987), No. 3, 493-496 (1988).
